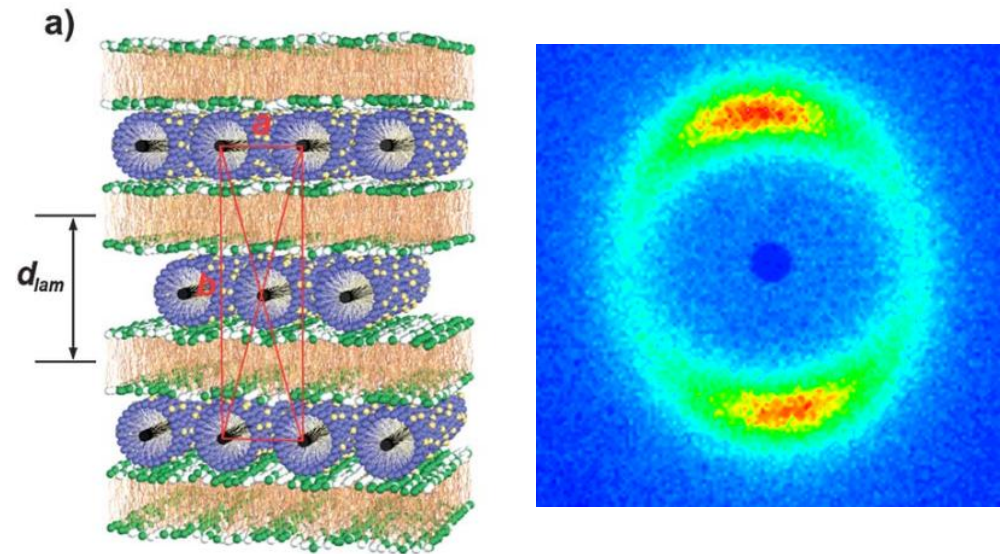
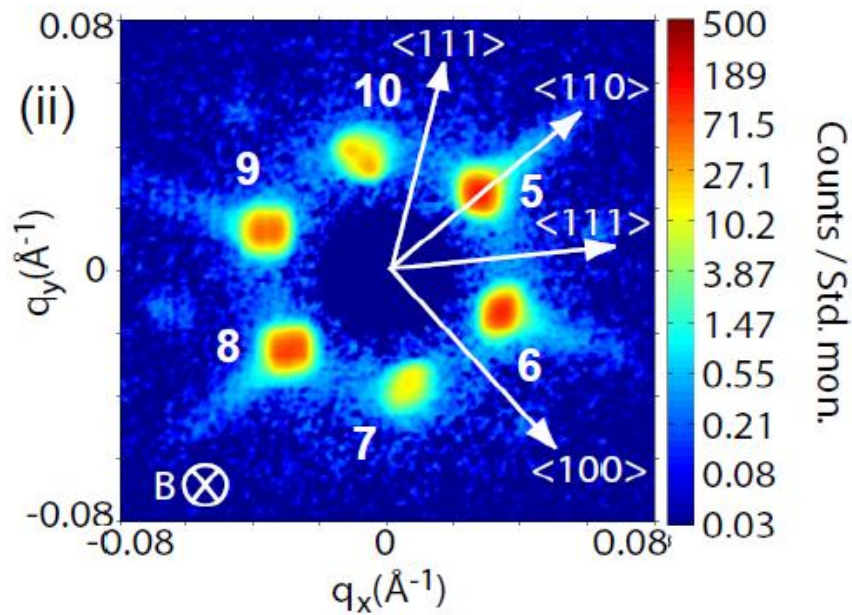


# Physics with Neutrons II, SS 2016



## Lecture 6, 6.6.2016

- VL1: Repetition of winter term, basic neutron scattering theory
- VL2: SANS, theory and applications
- VL3: Neutron optics
- VL4: Reflectometry and dynamical scattering theory
- VL5: Diffuse neutron scattering
- ➔ • VL6: Magnetic scattering cross section
- VL7: Magnetic structures and structure analysis
- VL8: Polarized neutrons and 3d-polarimetry
- VL9: Inelastic scattering on magnetism
- VL10: 4.7.2016 (8:30!!) Phase transitions and critical phenomena as seen by neutrons
- VL11: Spin echo spectroscopy

➔ Exam: Please register until 30.6.2016



Reminder:  
Diffuse Neutron Scattering:  
Looking between the Bragg spots

➔ Master formula for elastic coherent scattering on crystals

$$\frac{d\sigma}{d\Omega_{coh.el}} = N_0 \frac{(2\pi)^3}{v_0} e^{-2W(\vec{k})} \sum_{\vec{\tau}} |S_{\vec{\tau}}|^2 \delta(\vec{k} - \vec{\tau})$$

Normalization

Debye-Waller

Structure factor

Bragg positions

## Definition:

Diffuse scattering is scattering that arises from departures of any kind from a perfect regular crystal lattice

- Phonons: Thermal diffuse scattering (TDS)
- Distortions due to defects: Huang scattering
- Static or dynamic displacements
- Substitution or interstitial defects
- Stacking faults
- (Micro-)Domains
- Magnetic effects
- .....

➔ Diffuse neutron scattering can be static or dynamic

What we already know: Incoherent inelastic scattering

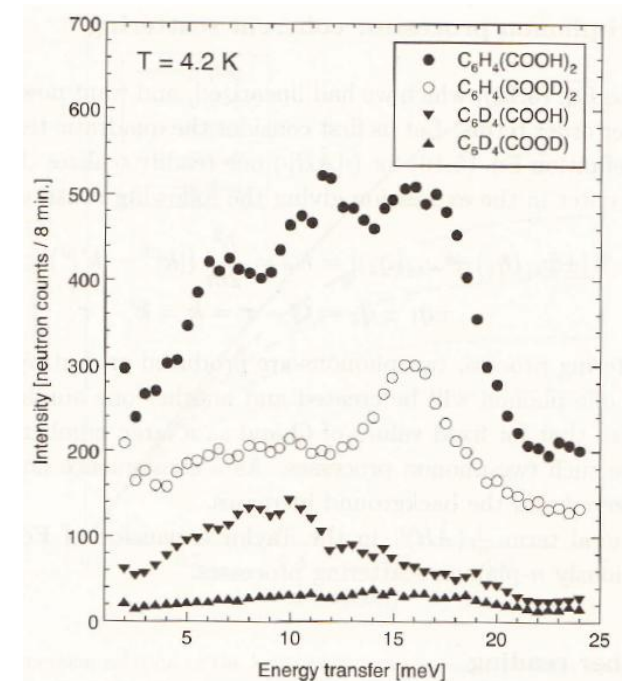
➡ Flat in Q (apart from DW)

➡ Gives Phonon DOS

$$\frac{d^2\sigma}{d\Omega d\omega} = \frac{1}{4M} \frac{k'}{k} (\langle b^2 \rangle - \langle b \rangle^2) e^{-W(\mathbf{Q})} \times \langle (\mathbf{Q} \cdot \mathbf{e}_s(\mathbf{q}))^2 \rangle \cdot \frac{g(\omega)}{\omega} \cdot \left[ \coth \frac{\hbar\omega}{2k_B T} \pm 1 \right]$$

With phonon DOS  $g(\omega)$

$$\int_0^\infty g(\omega) d\omega = 3N$$

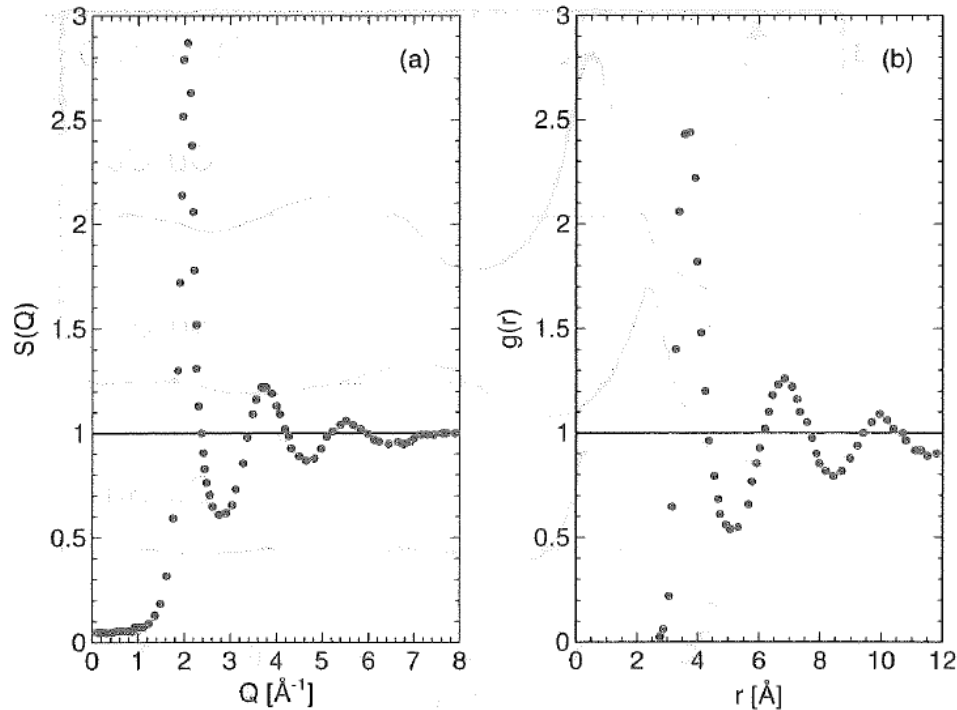


What we already know: Scattering on liquids and amorphous materials

➡ Measure of pair correlation function  $g(r)$

➡ Quasielastic scattering: Diffusion

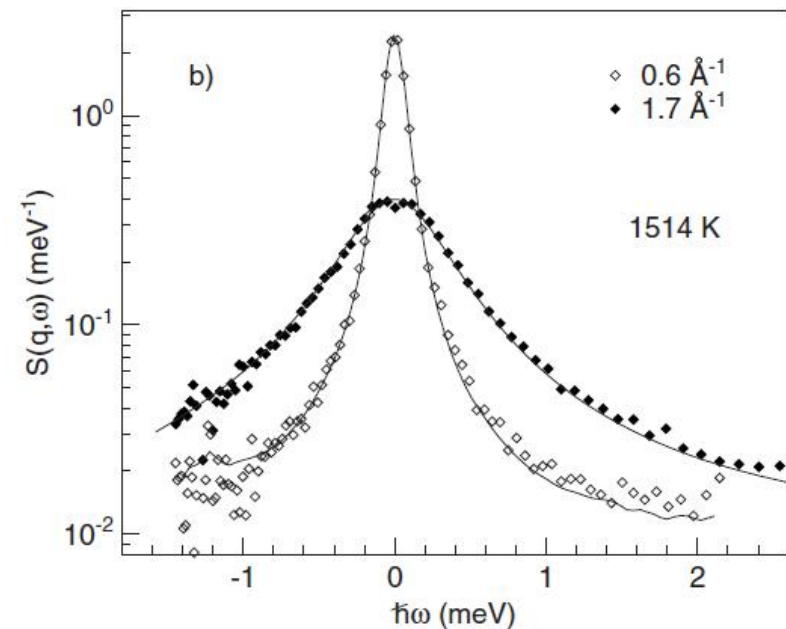
Static structure factor:  
Deviations of the mean density  $n(r)$



Static structure factor:  
Scattering function

$g(r)$  pair corr. function:  
Deviations from mean density  $n(r)$ .

Dynamic structure factor:  
Diffusive processes

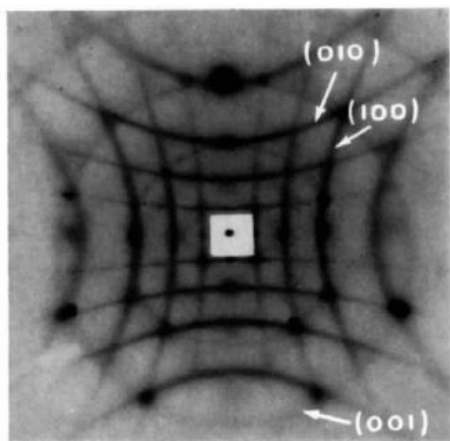


$S(Q, \omega) \xleftrightarrow{\text{FT}} G(r, t)$   
QENS, peaked at  $\omega=0$

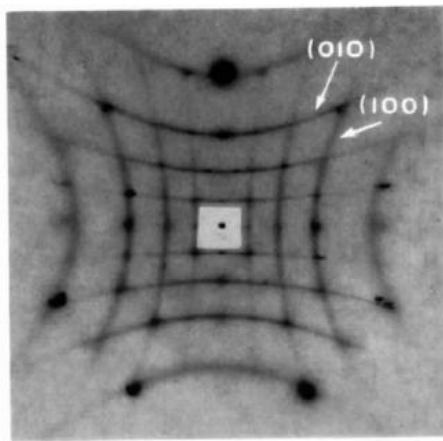


## Example 1: Ferroelectric Perovskite $\text{KNbO}_3$

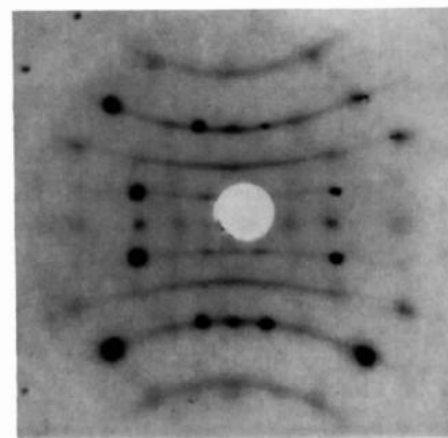
Here: Diffuse X-ray scattering (1970s)



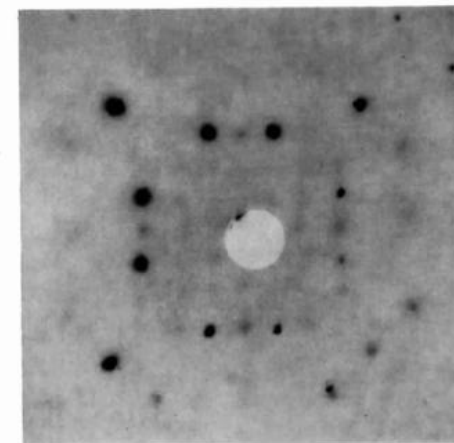
(a)



(b)



(c)



(d)

Two alternative models for diffuse sheets:

➔ Formation of linear chains of correlated vibrational displacements due to a specific soft mode?  
Thermal diffuse scattering?

➔ Formation of linear chains of correlated local displacements?

Ferroelectric Perovskite  $\text{KNbO}_3$

Compare to  $\text{PbTiO}_3$  and  $\text{BaTiO}_3$

??

Formation of linear chains of correlated vibrational displacements due to a specific soft mode?  
Thermal diffuse scattering?

Formation of linear chains of correlated local displacements?

➔ Soft mode similar for  $\text{PbTiO}_3$ ,  $\text{BaTiO}_3$  and  $\text{KNbO}_3$

➔ Different local symmetry

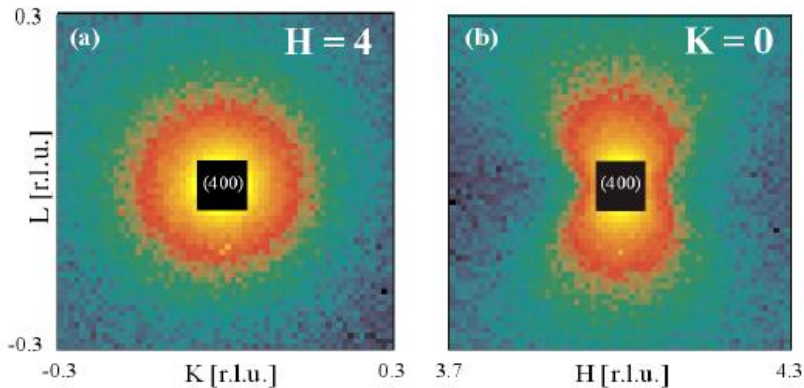
➔ Expect similar results

➔ No diffuse sheets for  $\text{PbTiO}_3$

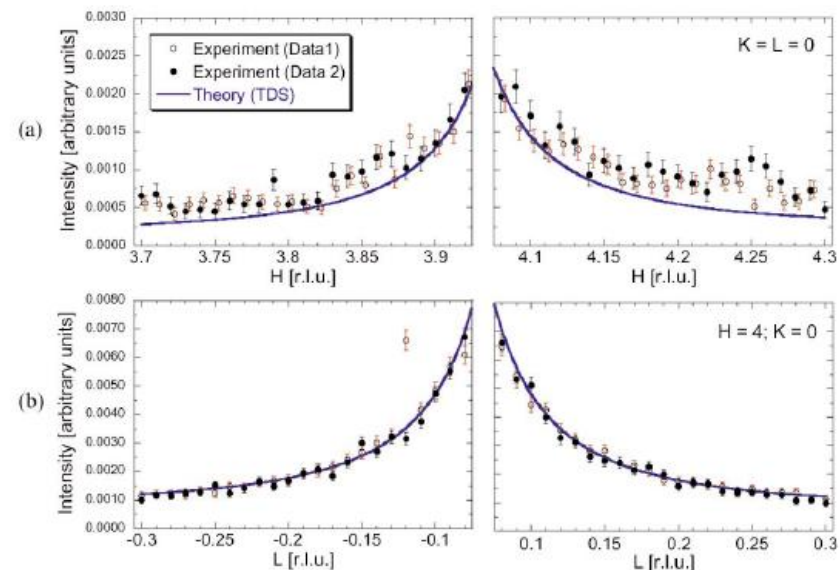
➔ Expect different result



Thermal diffuse scattering  $\text{PbTiO}_3$



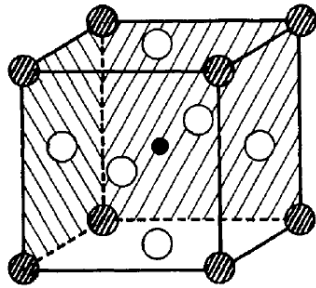
Diffuse intensity around Bragg positions





## Ferroelectric Perovskite $\text{KNbO}_3$

Formation of linear chains of correlated local displacements



● Ti ou Nb    ○ Oxygène    ▨ Ba ou K

Fig. 5. La structure pérovskite idéale.

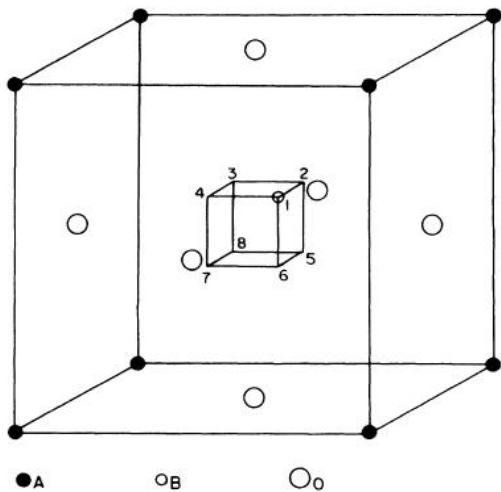
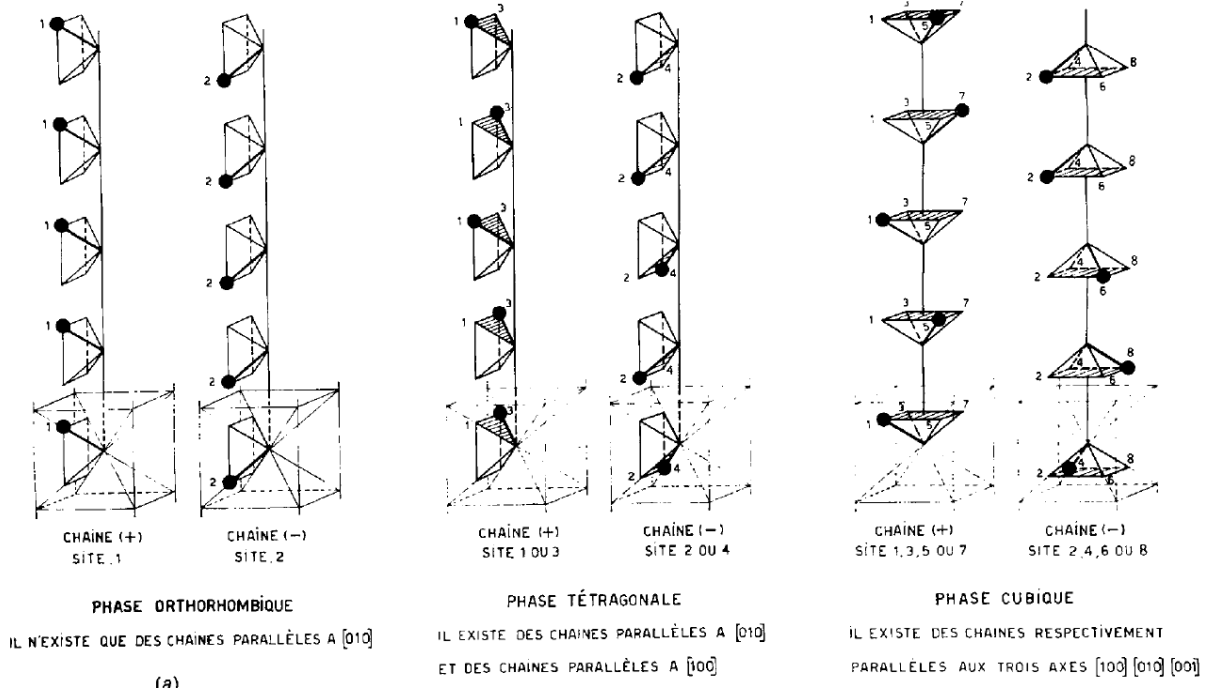
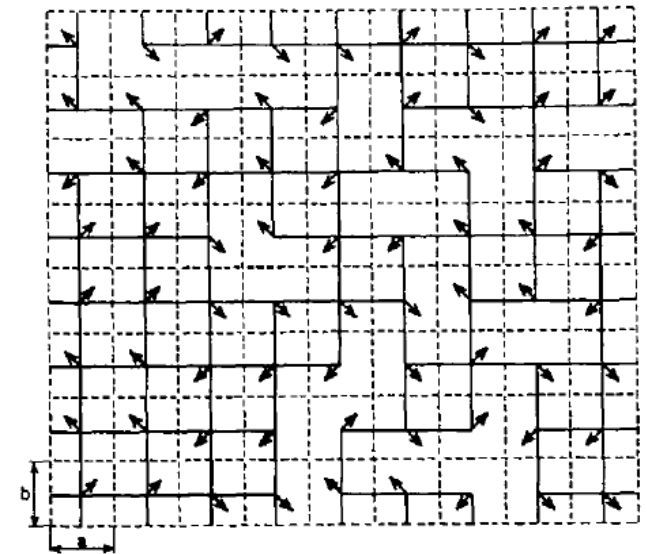


Fig. 10. Le croisement des chaînes de corrélation en projection sur un plan (001) dans la phase tétragonale. Les flèches schématisent le déplacement de l'atome central projeté sur le plan (001); les mailles appartenant à une même chaîne sont reliées en trait plein. En phase tétragonale la composante du déplacement perpendiculairement au plan de figure est constante; en phase cubique au contraire il existerait un système de chaînes analogue suivant [001].



## Example 2: Simple binary alloy on a bcc lattice

Crystallographic site can be occupied by

- two specific atoms
- half occupancy
- random occupancy

Simple rule:

- Similar atoms with scattering length  $b$
- No imaginary part
- No variation of  $b$
- Exactly one atom per unit cell occupied, however, randomly distributed

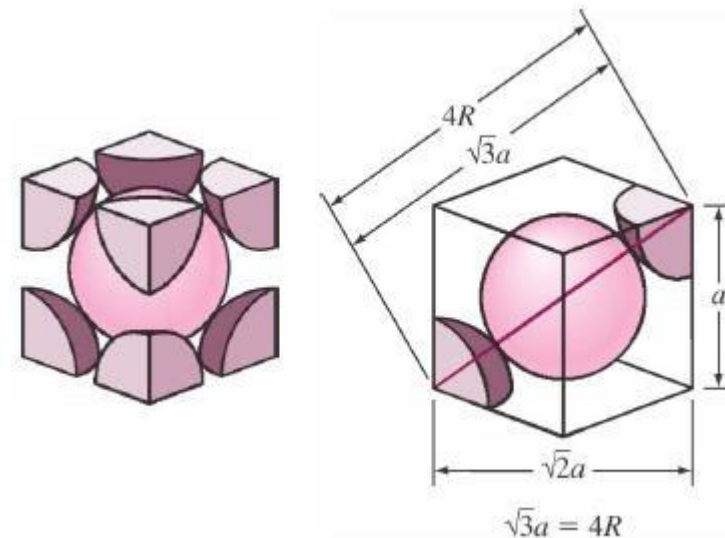
➔ „Random“ distribution, flat in  $Q$

Diffuse scattering

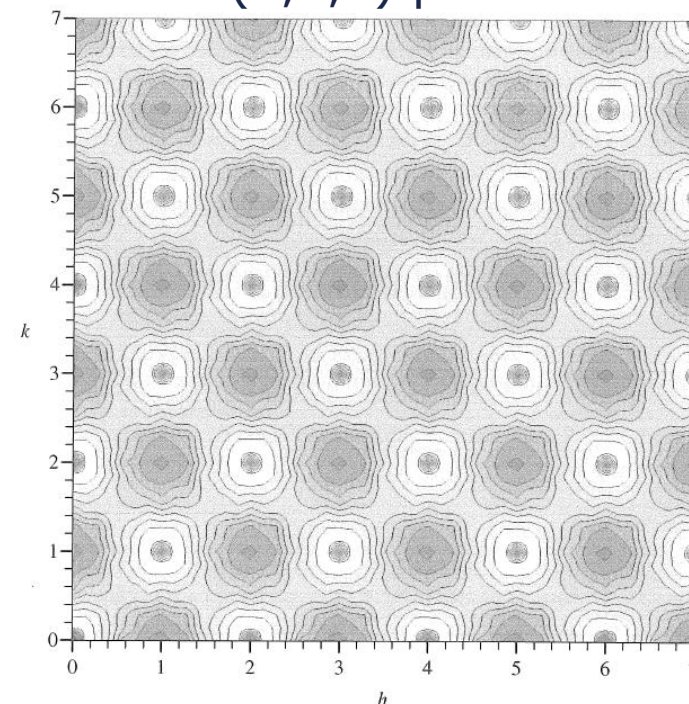
$$D^2 = \left| \sum_u s_u i b \sin(\mathbf{Q} \cdot \mathbf{d}) \exp(i\mathbf{Q} \cdot \mathbf{R}_u) \right|^2$$

+ Bragg scattering

$$\left( \frac{d\sigma}{d\Omega} \right)_{\text{coh}} = b^2 \sum_{j,j'} \exp(i\mathbf{Q} \cdot (\mathbf{R}_j - \mathbf{R}_{j'}))$$



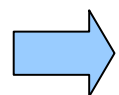
( $h, k, 0$ ) plane



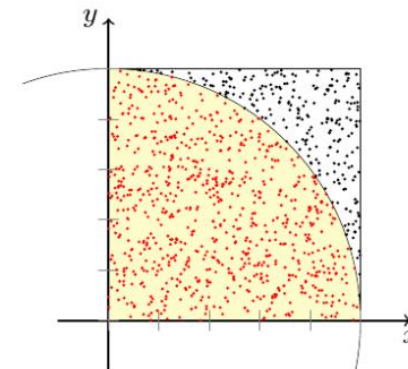
Powerful tool for modeling diffuse scattering

## Monte Carlo simulations (MC)

- 1) Choose appropriate descriptions of potentials
- 2) Generate arrangement of N atoms including boundary conditions
- 3) Calculate system energy
- 4) Randomly move atom(s)
- 5) Recalculate energy
- 6) Refine energy of the system with Monte Carlo method

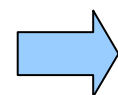


Working from the potential to the structure



## Reverse Monte Carlo simulations (RMC)

- 1) Generate arrangement of N atoms including boundary conditions and hard core potential to avoid overlap
- 2) Calculate  $F(Q)$  from arrangements
- 3) Randomly move atom(s)
- 4) Recalculate scattering signal  $F(Q)$
- 5) Refine  $F(Q)$  of the system with Monte Carlo method



Working from the scattering signal to the structure

Diffuse neutron scattering: Lets try a classification:

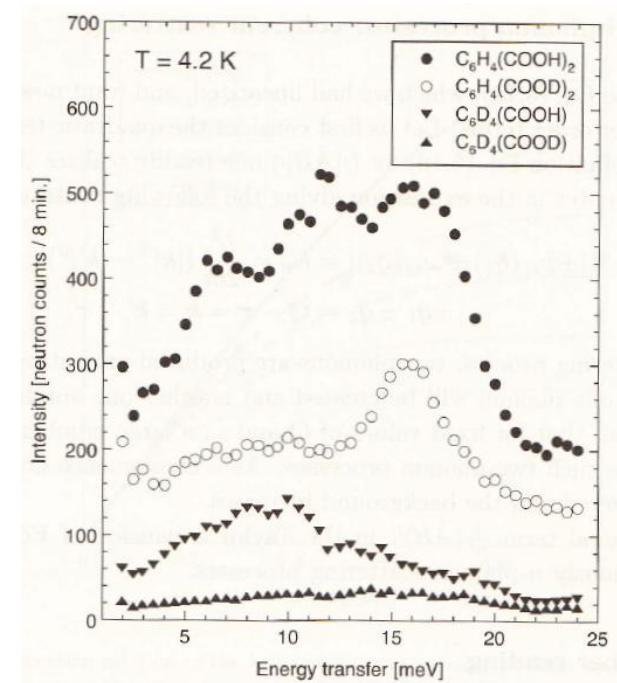
Diffuse background (flat in Q)

➔ Incoherent elastic scattering, random disorder (isotope, spin, voids....)

$$\frac{d\sigma}{d\Omega_{inc.}} = (\langle b^2 \rangle - \langle b \rangle^2) \sum_{j=j'} e^{-i\kappa(R_{j'} - R_j)} = N(\langle b^2 \rangle - \langle b \rangle^2)$$

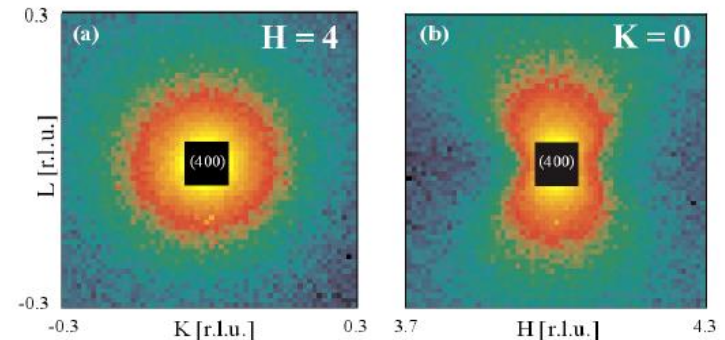
➔ Incoherent inelastic scattering, phonon DOS (Debye Waller, not really flat)

$$\frac{d^2\sigma}{d\Omega d\omega} = \frac{1}{4M} \frac{k'}{k} (\langle b^2 \rangle - \langle b \rangle^2) e^{-W(Q)} \times \langle (Q \cdot e_s(\mathbf{q}))^2 \rangle \cdot \frac{g(\omega)}{\omega} \cdot \left[ \coth \frac{\hbar\omega}{2k_B T} \pm 1 \right]$$



Odd Bragg peak shape (Butterflies, ellipsoids..)

➔ Typically thermal diffuse scattering, Einstein model for optical phonons, measurement of lattice elasticity





## Diffuse neutron scattering: Lets try a classification:

### Bragg rods and planes

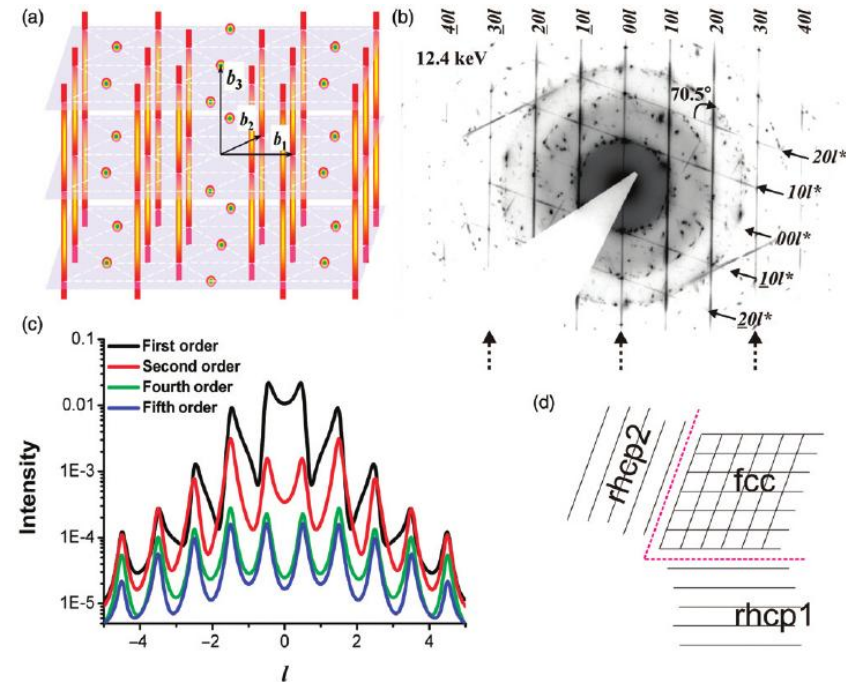
➔ Other than 3D order

Line order (1D)

➔ Bragg planes

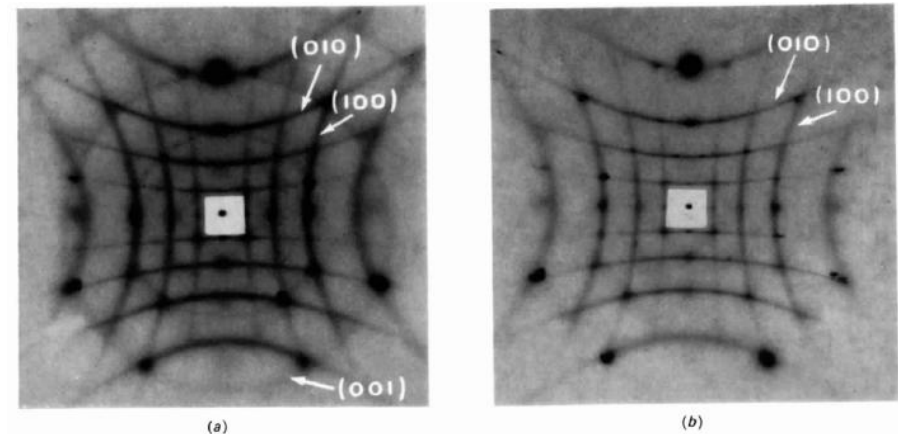
Planar correlations (2D order)

➔ Bragg rods



### Diffuse but with structure in Q

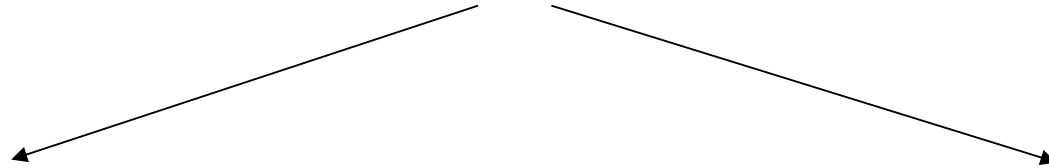
➔ Partial order /suborder of systems  
See examples 1 and 2,  
Liquid and amorphous samples





In general: Size of distortion  $\sim L$   
Spread in  $\Delta q \sim 2\pi/L$

Try to minimize parameter space



Short range features, spread  
over large Q

Work in real space  
(correlation approach,  
microdomain approach)

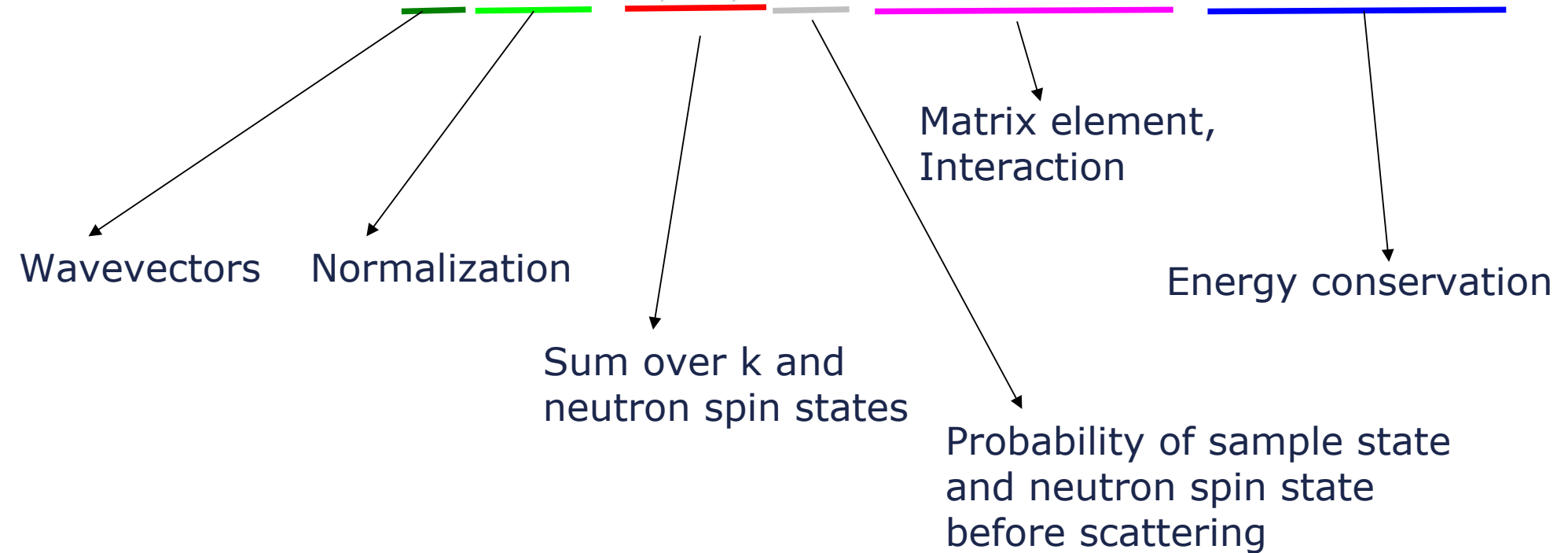
Long range features, spread  
over small Q

Work in reciprocal space  
(Modulation wave approach)

# Magnetic neutron scattering: Basic cross section

Starting point (similar to Winter term, now including the spin state of the neutron)

$$\left( \frac{d^2\sigma}{d\Omega dE'} \right) = \frac{k'}{k} \left( \frac{m}{2\pi\hbar} \right)^2 \sum_{k',\sigma'} \sum_{k,\sigma} p_\lambda p_\sigma |\langle k'\sigma'\lambda' | U | k\sigma\lambda \rangle|^2 \delta(\hbar\omega + E_\lambda - E_{\lambda'})$$



Until now: Only nuclear scattering  
Interaction: Fermi pseudopotential

$$V(r) = \frac{2\pi\hbar^2}{m} b\delta(r)$$

## Nuclear neutron scattering

Interaction: Fermi pseudopotential  $V(r) = \frac{2\pi\hbar^2}{m} b\delta(r)$

- ➡ Scalar function
- ➡ Point like (delta function)
- ➡ For an incoming plane wave: s-wave scattering
- ➡ Spherical symmetry
- ➡ FT of delta function is constant

Magnetic neutron scattering:

Interaction: Magnetic moment of neutron interacts with local magnetic field

$$\hat{U} = \hat{\mu} \cdot \vec{H} \qquad \hat{U} = \hat{\mu} \cdot \vec{H} = \gamma \mu_N \hat{\sigma} \cdot \vec{H}$$

$$\vec{H} = \nabla \times \left( \frac{\vec{\mu}_e \times \vec{R}}{|\vec{R}|^3} \right) - \frac{e \vec{v}_e \times \vec{R}}{c |\vec{R}|^3}$$

Spin of electrons

Orbital momentum

- ➡ Vector (dipole-dipole) interaction
- ➡ Extended range, not point like
- ➡ No s-wave scattering
- ➡ No spherical symmetry
- ➡ Magnetic moments: Unpaired electrons



## Leaving away the maths

- ➡ For spin only scattering (neglect orbital momentum)
- ➡ For unpolarized neutrons (average over polarization states)
- ➡ For identical magnetic ions with localized moments

Master formula

$$\left( \frac{d^2\sigma}{d\Omega dE'} \right) = (\gamma r_0)^2 \frac{k'}{k} F^2(\vec{Q}) e^{-2W(\vec{Q})} \sum_{\alpha, \beta} \left( \delta_{\alpha, \beta} - \frac{Q_\alpha Q_\beta}{Q^2} \right) S^{\alpha\beta}(Q, \omega)$$

$$S^{\alpha\beta}(Q, \omega) = \sum_{j, j'} e^{i\vec{Q}(R_j - R_{j'})} \sum_{\lambda, \lambda'} p_\lambda \langle \lambda | S_{j'}^\alpha | \lambda' \rangle \langle \lambda | S_j^\beta | \lambda' \rangle \delta(\hbar\omega + E_\lambda - E_{\lambda'})$$

## Master formula

$$\left( \frac{d^2\sigma}{d\Omega dE'} \right) = (\gamma r_0)^2 \frac{k'}{k} F^2(\vec{Q}) e^{-2W(\vec{Q})} \sum_{\alpha, \beta} \left( \delta_{\alpha, \beta} - \frac{Q_\alpha Q_\beta}{Q^2} \right) S^{\alpha\beta}(Q, \omega)$$

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Dipole-dipole interaction: Angular dependence

Neutron polarization factor  $\left( \delta_{\alpha, \beta} - \frac{Q_\alpha Q_\beta}{Q^2} \right)$

➡ Only moments perpendicular to Q contribute to magnetic scattering  
Don't confuse with polarized neutrons!

## Master formula

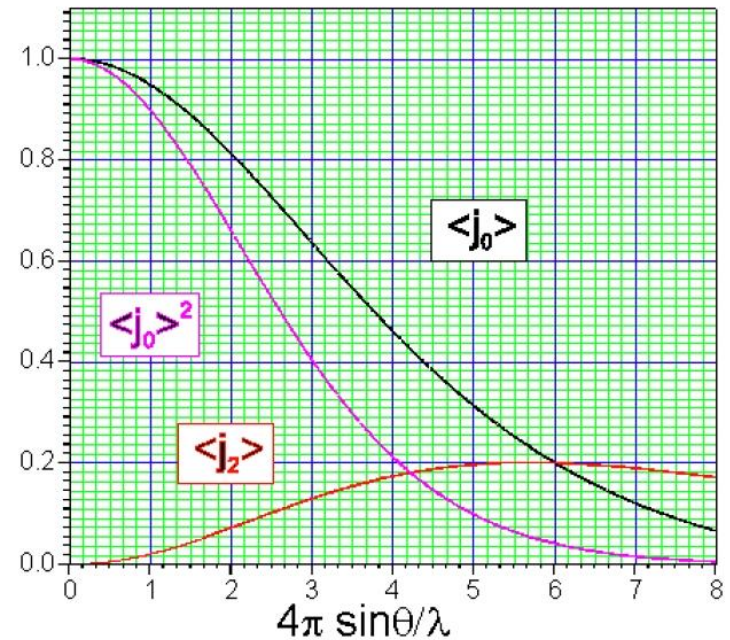
$$\left( \frac{d^2\sigma}{d\Omega dE'} \right) = (\gamma r_0)^2 \frac{k'}{k} \underline{F^2(\vec{Q})} e^{-2W(\vec{Q})} \sum_{\alpha, \beta} \left( \delta_{\alpha, \beta} - \frac{Q_\alpha Q_\beta}{Q^2} \right) S^{\alpha\beta}(Q, \omega)$$

$$S^{\alpha\beta}(Q, \omega) = \sum_{j, j'} e^{i\vec{Q}(R_j - R_{j'})} \sum_{\lambda, \lambda'} p_\lambda \langle \lambda | S_{j'}^\alpha | \lambda' \rangle \langle \lambda | S_j^\beta | \lambda' \rangle \delta(\hbar\omega + E_\lambda - E_{\lambda'})$$

## Dipole-dipole interaction: Magnetic form factor

- ➡ Fourier transform of electron cloud
- ➡ Useful to discriminate magnetic/nuclear scattering
- ➡ Check the tables for each ion!

**Fe<sup>3+</sup>: 3d<sup>5</sup> 6S**



For the case of orbital momentum + spin

$$\mu = -\mu_b(L + 2S)$$

Effective angular  
momentum operator

$$\hat{S}_j^\alpha = \frac{1}{2}g\hat{J}_j^\alpha$$

Landé splitting factor

$$g = 1 + \frac{J(J + 1) - L(L + 1) + S(S + 1)}{2J(J + 1)}$$

⇒ Approximation for small Q, spin+orbital momentum

Magnetic scattering function

$$S^{\alpha\beta}(Q, \omega) = \sum_{j,j'} e^{i\vec{Q}\cdot(R_j - R_{j'})} \langle \underline{S_{j'}^\alpha(0) S_{j'}^\beta(t)} \rangle e^{-i\omega t} dt$$

Spin correlation function:

Correlation of magnetic moment at site j, time t=0  
and site j', time t=t

⇒ Fourier transform measured with neutrons!

Why neutrons are so unique:

Fluctuation dissipation theorem

$$S^{\alpha\beta}(Q, \omega) = \frac{N\hbar}{\pi} (1 - e^{-\frac{\hbar\omega}{k_b T}})^{-1} \text{Im} \chi^{\alpha\beta}(\vec{Q}, \omega)$$

Neutrons directly measure generalized susceptibility tensor

$$M^\alpha(\vec{Q}, \omega) = \chi^{\alpha\beta}(\vec{Q}, \omega) H^\beta(\vec{Q}, \omega)$$