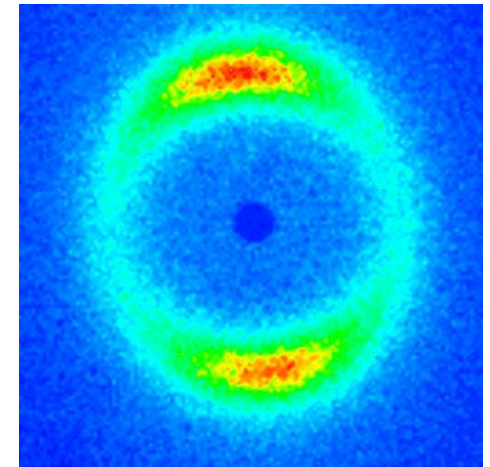
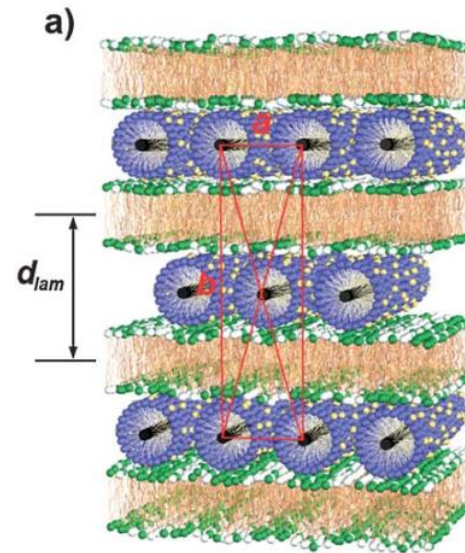
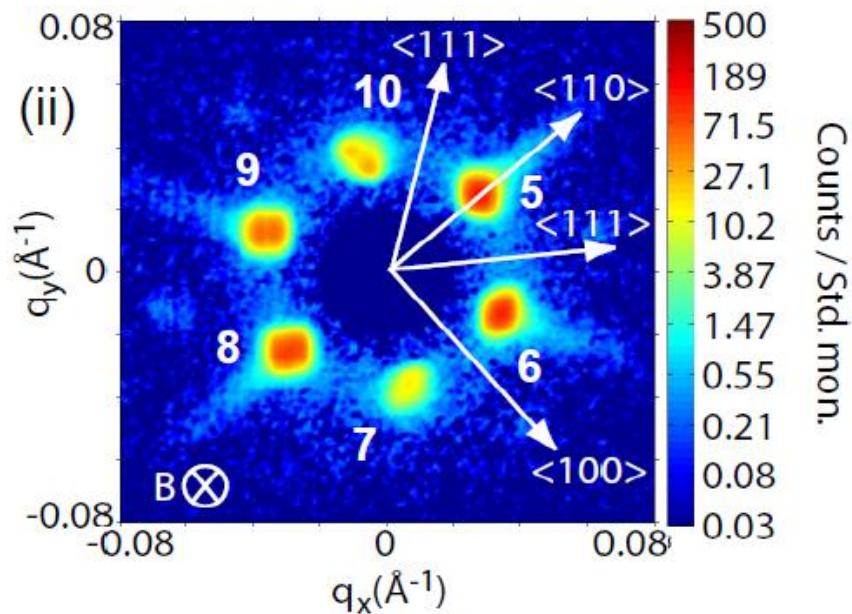


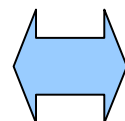
# Physics with Neutrons II, SS 2016



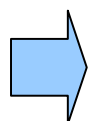
## Lecture 2, 25.4.2016

Pair correlation function  $G(r,t)$  useful for description of liquids

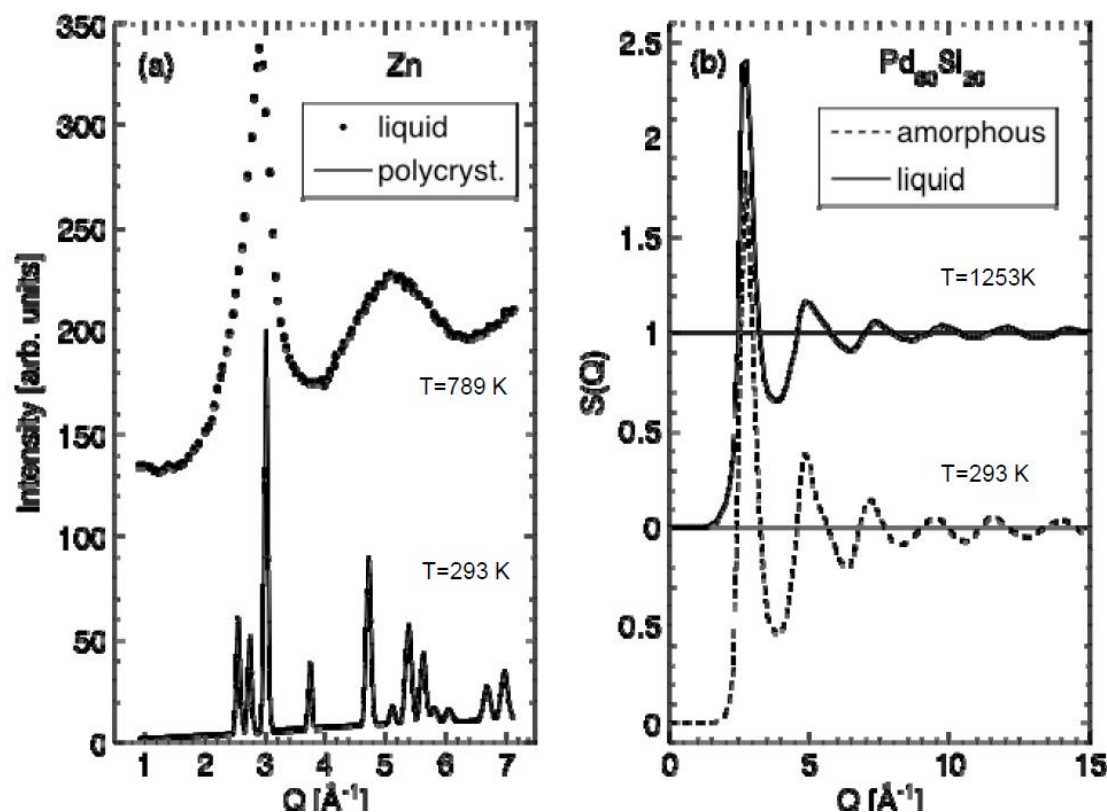
Liquid (amorphous) sample



Crystalline sample



Similar density, no LRO, only short range correlations



Static structure factor: Looking at deviations of the mean density  $n(\mathbf{r})$

$$G'(\mathbf{r}) = \frac{1}{N} \int \langle n(\mathbf{r}' - \mathbf{r}) - \langle n(\mathbf{r}' - \mathbf{r}) \rangle \rangle (n(\mathbf{r}') - \langle n(\mathbf{r}') \rangle) \mathrm{d}\mathbf{r}'$$

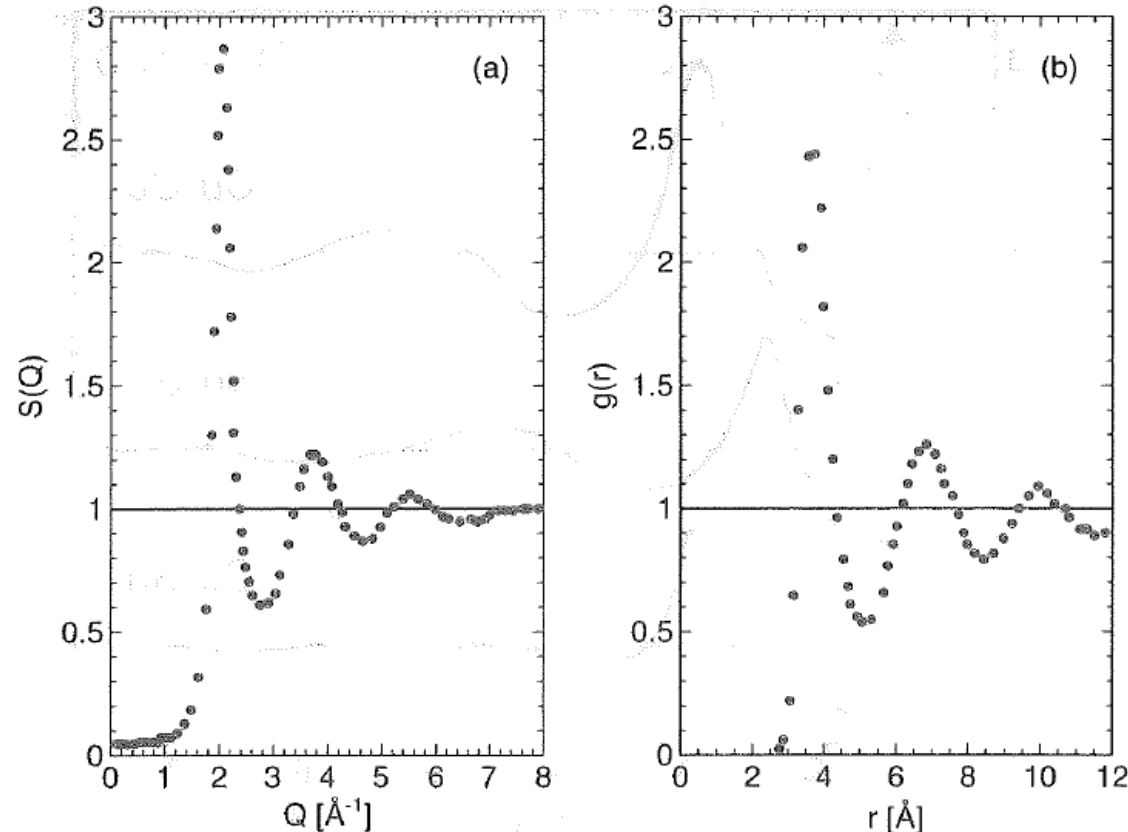
Elastic scattering from liquids

$$\frac{\mathrm{d}\sigma}{\mathrm{d}\Omega} = N \langle b \rangle^2 \left( 1 + \int (g(\mathbf{r}) - n_0) e^{i\mathbf{Q}\mathbf{r}} \mathrm{d}\mathbf{r} \right)$$

$g(\mathbf{r})$ : pair correlation function

$$S(Q) = 1 + 4\pi \int_0^\infty (g(r) - n_0) \frac{\sin Qr}{Qr} r^2 \mathrm{d}r$$

Static structure factor: Looking at deviations of the mean density  $n(r)$



Static structure factor:  
Scattering function

$g(r)$  pair correlation function:  
Deviations from mean density  $n(r)$ .

Limit  $Q \rightarrow \infty$      $S(Q=\infty)=1$

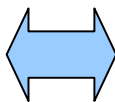
Limit  $Q \rightarrow 0$      $S(Q \rightarrow 0) = n_0 \kappa_T k_B T$     Isothermal compressibility

# SANS – applications and extensions

# SANS – Basic theory

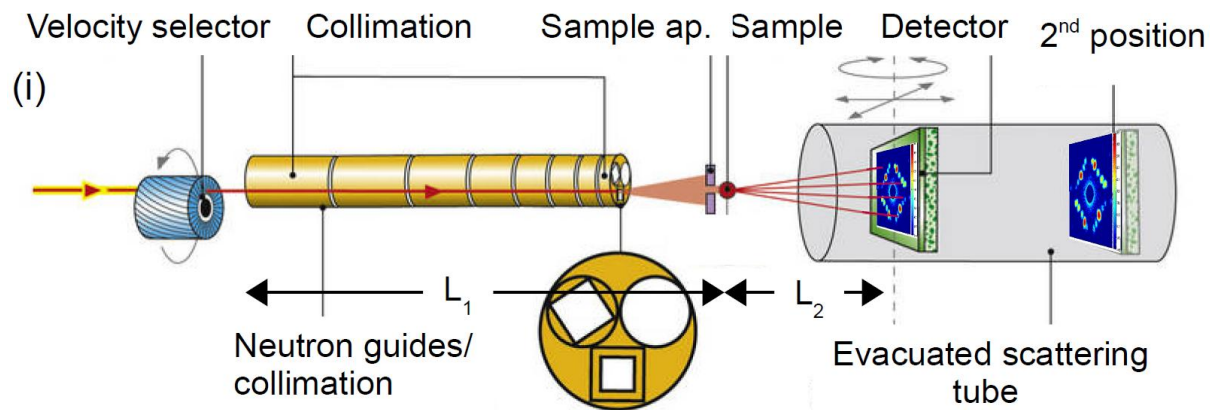


Large scales in real space  
 $10\text{-}40000\text{\AA}$



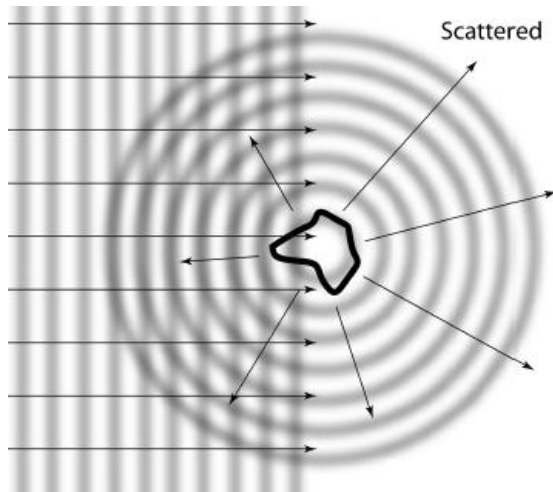
Low  $Q$ , small scattering angles  
 $0.54\text{\AA}^{-1} - 6 \cdot 10^{-4}\text{\AA}^{-1}$

Diffractometer specialized for small scattering angles

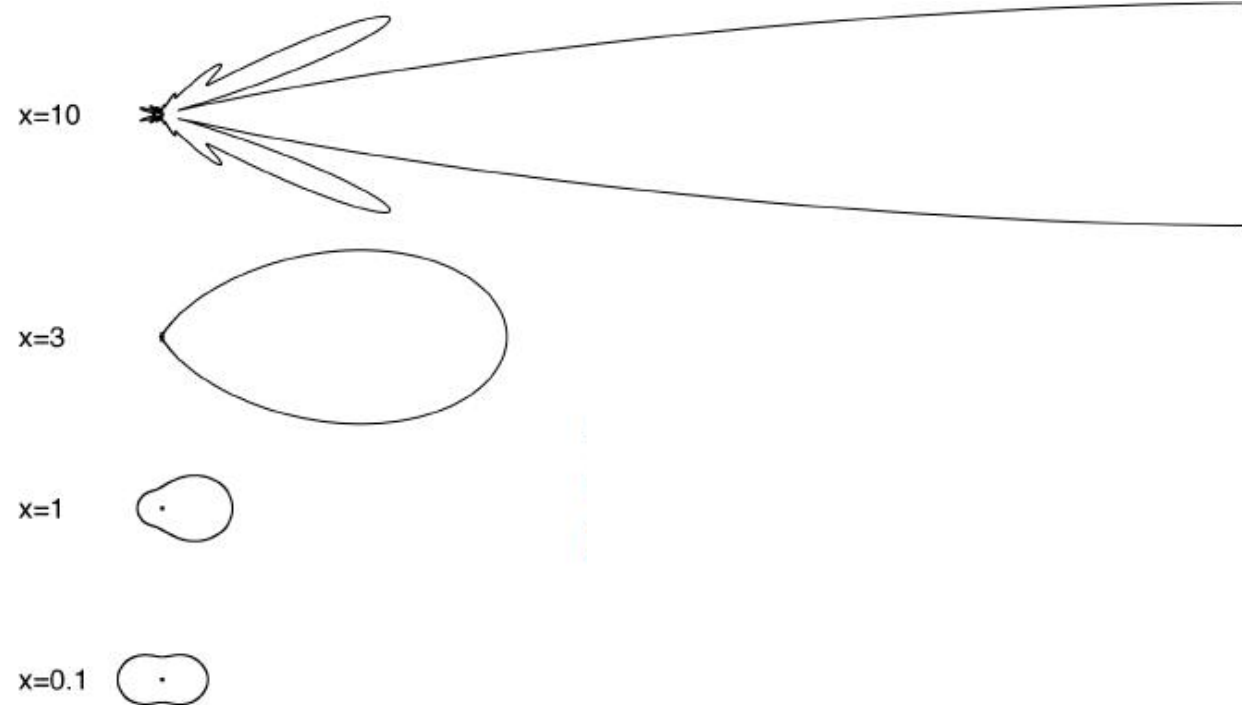


# Forward scattering-Why is it useful?

Backscattering      Forward scattering



Phase function for different  $x$



$$x = \frac{2\pi r}{\lambda}$$

➡ Lots of intensity scattered in forward direction



Starting point: coherent elastic cross section

$$\frac{d\sigma}{d\Omega} = \int_{-\infty}^{\infty} \frac{d^2\sigma}{d\Omega dE'} d(\hbar\omega) = \sum_{j,j'} b_j b_{j'} \langle e^{-i\mathbf{q}\cdot\mathbf{r}_{j'}} e^{-i\mathbf{q}\cdot\mathbf{r}_j} \rangle$$

Sum over identical atoms

$$\frac{d\sigma}{d\Omega}(\mathbf{q}) = \frac{1}{N} \left| \sum_i^N b_i e^{i\mathbf{q}\cdot\mathbf{r}_i} \right|^2$$

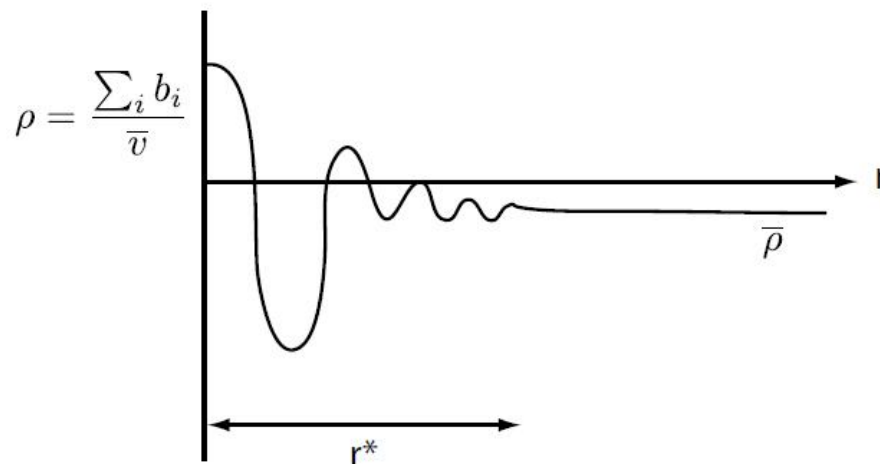
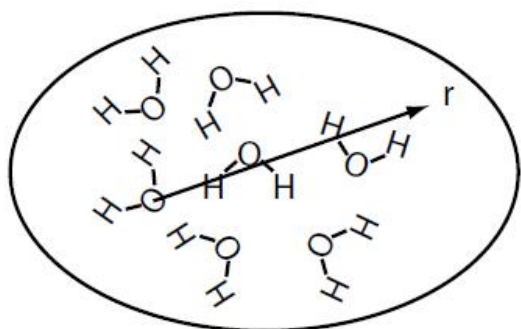
➡ SANS: low  $q$  averages over large  $r$

Scattering length  $b$  ➡ Scattering length density

$$\rho = \frac{\sum_i^n b_i}{V}$$

$$\rho(\mathbf{r}) = b_i \delta(\mathbf{r} - \mathbf{r}_i)$$

Example: water



Insert scattering length density into coherent elastic cross section:

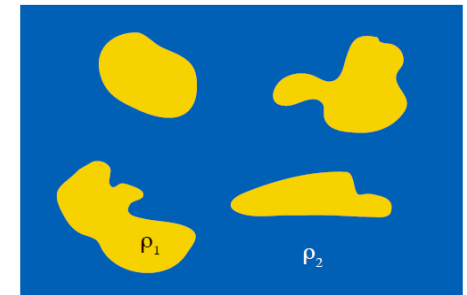
Rayleigh-Gans Equation 
$$\frac{d\Sigma}{d\Omega}(\mathbf{q}) = \frac{N}{V} \frac{d\sigma}{d\Omega}(\mathbf{q}) = \frac{1}{V} \left| \int_V \rho(\mathbf{r}) e^{i\mathbf{q} \cdot \mathbf{r}} d\mathbf{r} \right|^2$$

➡ SANS measures inhomogeneities of scattering length density

Assume a general two phase system

$$V = V_1 + V_2$$

$$\rho(r) = \begin{cases} \rho_1 & \text{in } V_1 \\ \rho_2 & \text{in } V_2 \end{cases}$$



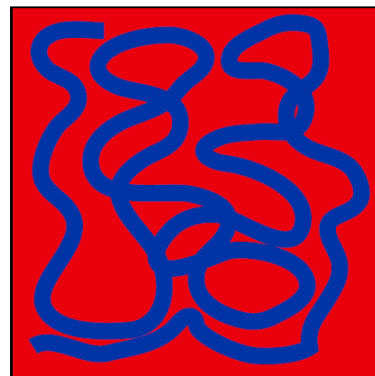
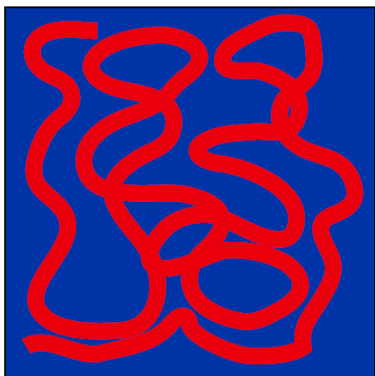
Split up the  
integral over  
the sample,  
break up into  
the two  
subvolumes

$$\frac{d\Sigma}{d\Omega}(\mathbf{q}) = \frac{1}{V} \left| \int_{V_1} \rho_1 e^{i\mathbf{q} \cdot \mathbf{r}} d\mathbf{r}_1 + \int_{V_2} \rho_2 e^{i\mathbf{q} \cdot \mathbf{r}} d\mathbf{r}_2 \right|^2$$

$$\frac{d\Sigma}{d\Omega}(\mathbf{q}) = \frac{1}{V} \left| \int_{V_1} \rho_1 e^{i\mathbf{q} \cdot \mathbf{r}} d\mathbf{r}_1 + \rho_2 \left\{ \int_V e^{i\mathbf{q} \cdot \mathbf{r}} d\mathbf{r} - \int_{V_1} e^{i\mathbf{q} \cdot \mathbf{r}} d\mathbf{r}_1 \right\} \right|^2$$

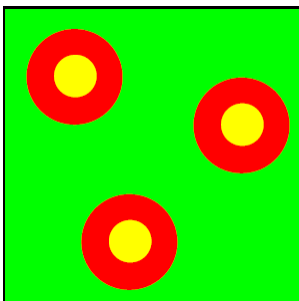
➡ 
$$\frac{d\Sigma}{d\Omega}(\mathbf{q}) = \frac{1}{V} (\rho_1 - \rho_2)^2 \left| \int_{V_1} e^{i\mathbf{q} \cdot \mathbf{r}} d\mathbf{r}_1 \right|^2$$

$$\Rightarrow \frac{d\Sigma}{d\Omega}(\mathbf{q}) = \frac{1}{V}(\rho_1 - \rho_2)^2 \left| \int_{V_1} e^{i\mathbf{q} \cdot \mathbf{r}} d\mathbf{r}_1 \right|^2$$

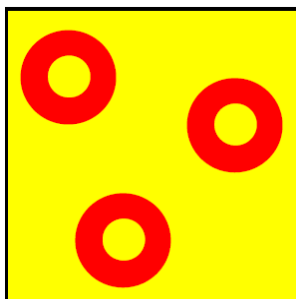


$\Rightarrow$  Principle of Babinet: Same coherent scattering

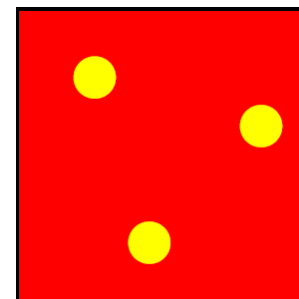
Contrast variation and contrast matching



Natural contrast



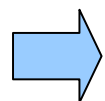
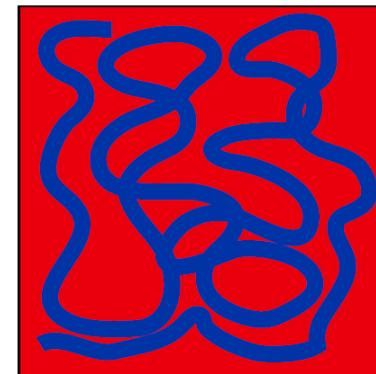
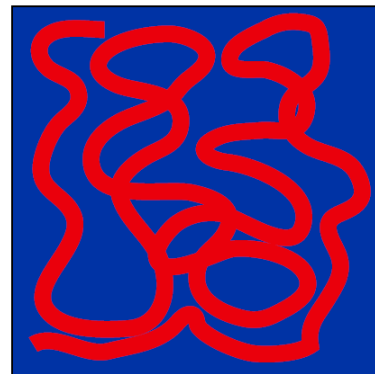
Shell visible



Core visible

$\Rightarrow$  Often: Mixture of H<sub>2</sub>O/D<sub>2</sub>O, isotope variation

$$\frac{d\Sigma}{d\Omega}(\mathbf{q}) = \frac{1}{V}(\rho_1 - \rho_2)^2 \left| \int_{V_1} e^{i\mathbf{q} \cdot \mathbf{r}} d\mathbf{r}_1 \right|^2$$



Split up the integral over the sample

$$\frac{d\Sigma}{d\Omega}(q) = \frac{N}{V}(\rho_1 - \rho_2)^2 V_p^2 P(q) S(q)$$

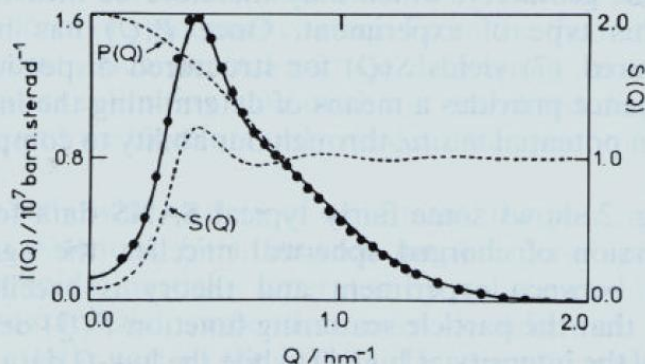


Fig. 2. Observed (●) and calculated (—) scattered intensity  $I(Q)$  as a function of momentum transfer  $Q$  for a charged micellar dispersion:  $0.03 \text{ mol dm}^{-3}$  hexadecyltrimethylammonium chloride in  $\text{D}_2\text{O}$  at 313 K. The functions  $P(Q)$  and  $S(Q)$  are discussed in the text. ( $1 \text{ barn sterad}^{-1} = 10^{-28} \text{ M}^2 \text{ sterad}^{-1}$ ).

Form factor  $P(Q)$

Interference of neutrons scattered at the same object

Shape, surface and density distribution of objects

Structure factor  $S(Q)$

Interference of neutrons scattered from different objects

Arrangement or superstructure of objects

Remember: Bragg scattering and Debye Waller

Form factor  $P(Q)$

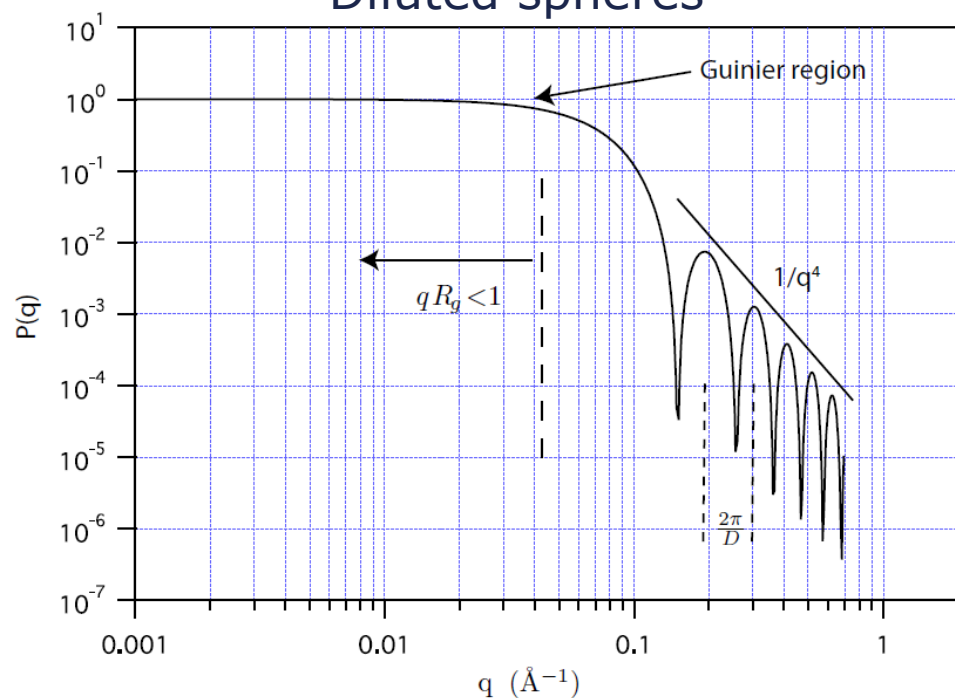
Interference of neutrons  
scattered at the same  
object

For isotropic solutions:

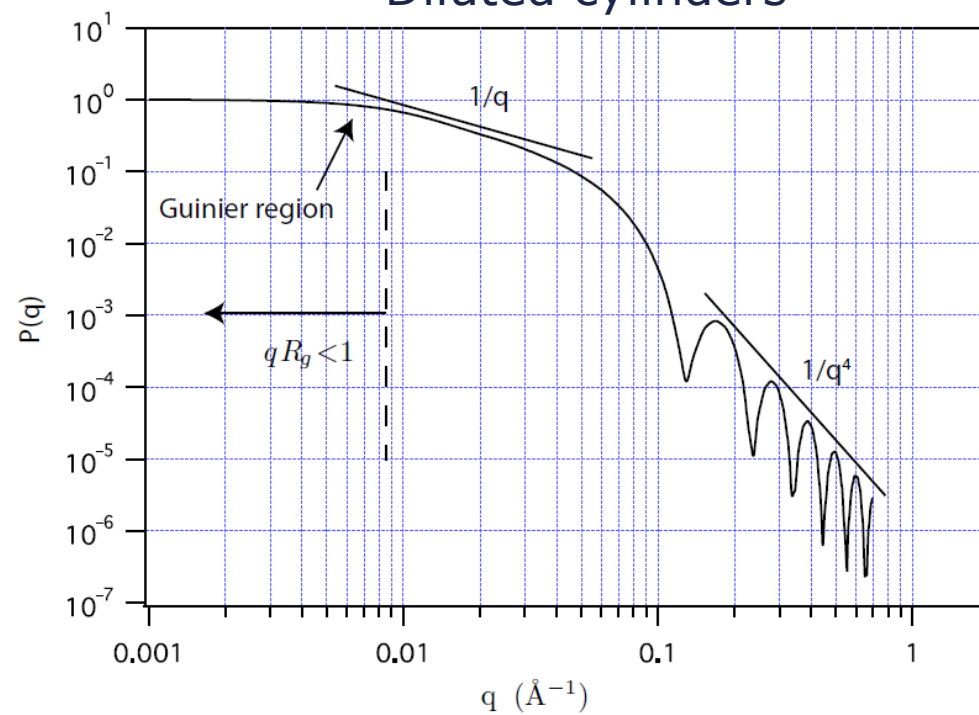
$$S(q) = 1 + 4\pi N_p \int_0^\infty [g(r) - 1] \frac{\sin(qr)}{qr} r^2 dr$$

$g(r)$  pair correlation function

Diluted spheres

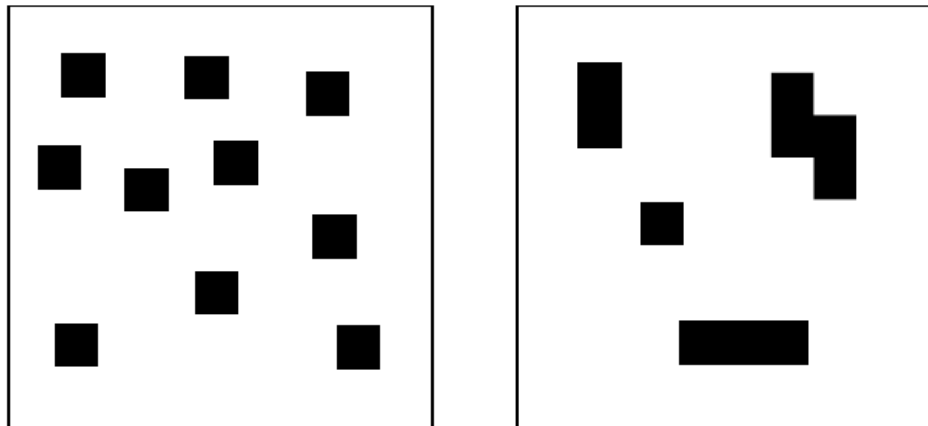


Diluted cylinders



More about that in the tutorial!

Two samples with 10% white and 90% black



Integrate with respect to  $Q$

$$\frac{d\Sigma}{d\Omega}(\mathbf{q}) = \frac{1}{V}(\rho_1 - \rho_2)^2 \left| \int_{V_1} e^{i\mathbf{q}\cdot\mathbf{r}} d\mathbf{r}_1 \right|^2$$

For a two phase system

$$\frac{Q}{4\pi} = Q^* = 2\pi^2 \phi_1 (1 - \phi_1) (\rho_2 - \rho_1)^2$$

Scattering invariant, total small angle scattering is constant!



# Porod and Guinier regime

Porod scattering for smooth surfaces and  
 $Q \gg 1/D$

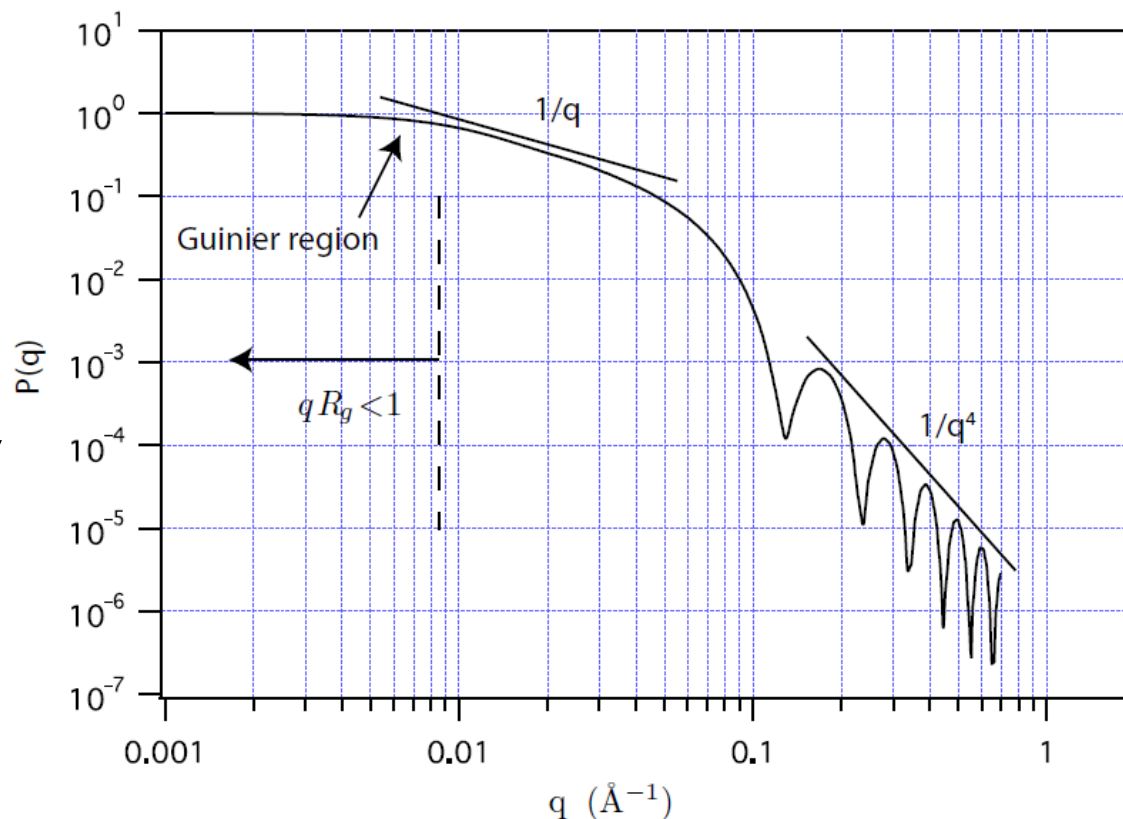
$$I(q) \propto (q)^{-4}$$

$$\frac{\pi}{Q^*} \cdot \lim_{q \rightarrow \infty} (I(q) \cdot q^4) = \frac{S}{V}$$

Guinier scattering for dilute, monodisperse and isotropic solutions of particles:  
 $QR_G \ll 1$

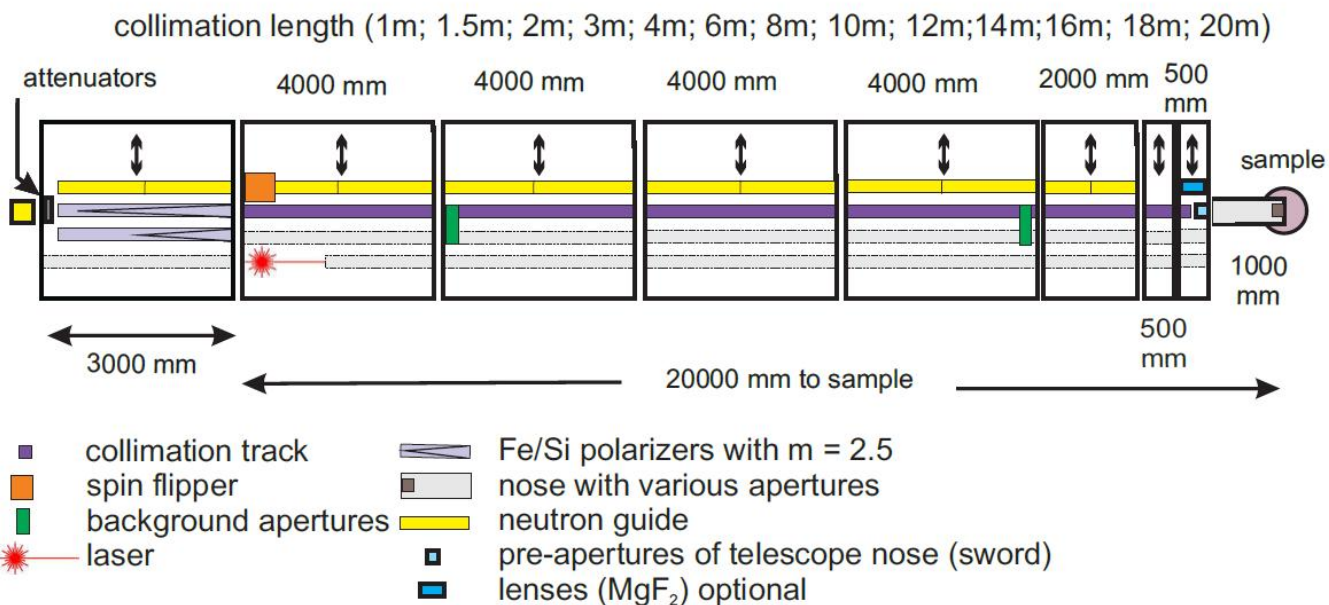
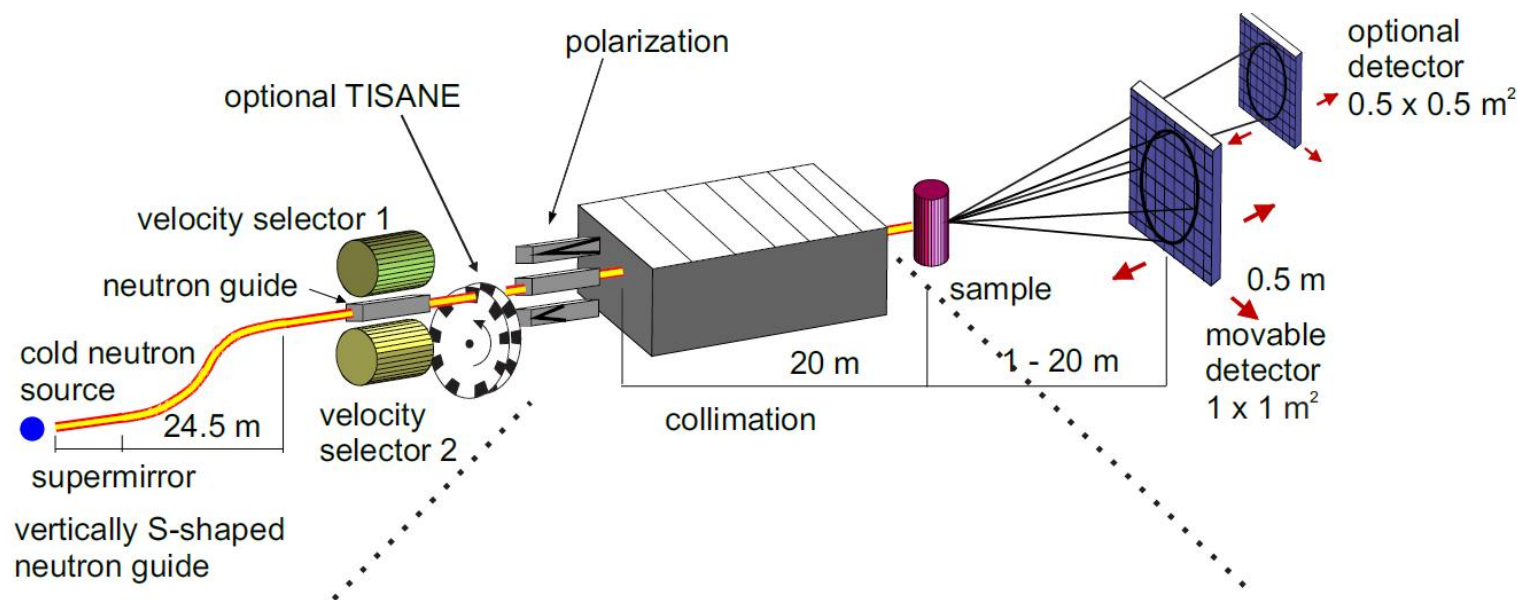
$$I(q) = I(0)e^{\frac{-(qR_g)^2}{3}}$$

Form factor for diluted cylinders  
radius 30Å, length 400Å  
No structure factor!



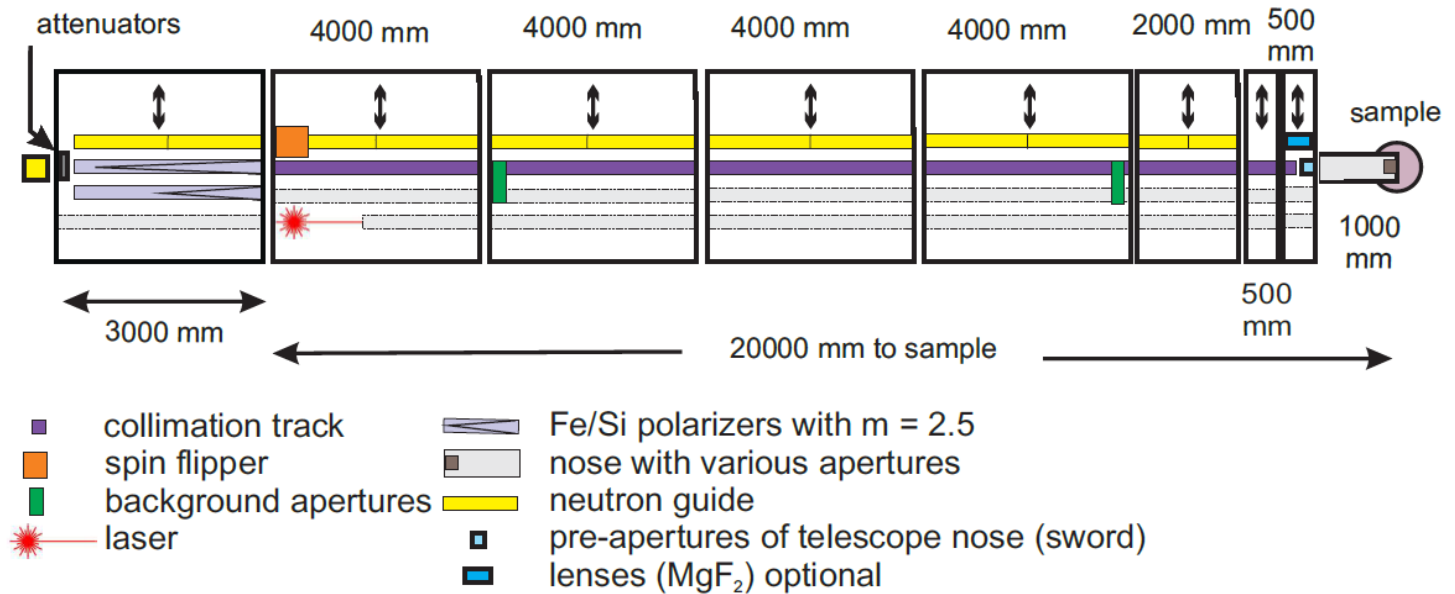
# SANS – Instrumentation

# Elements of a SANS Instrument





Velocity selector



Collimation: Define resolution and intensity

Aperture system/neutron guides (supermirror)

Alignment extremely critical

Well-defined and homogenous wavelength /divergence profiles

Transmission polarizer for the use of polarized neutrons

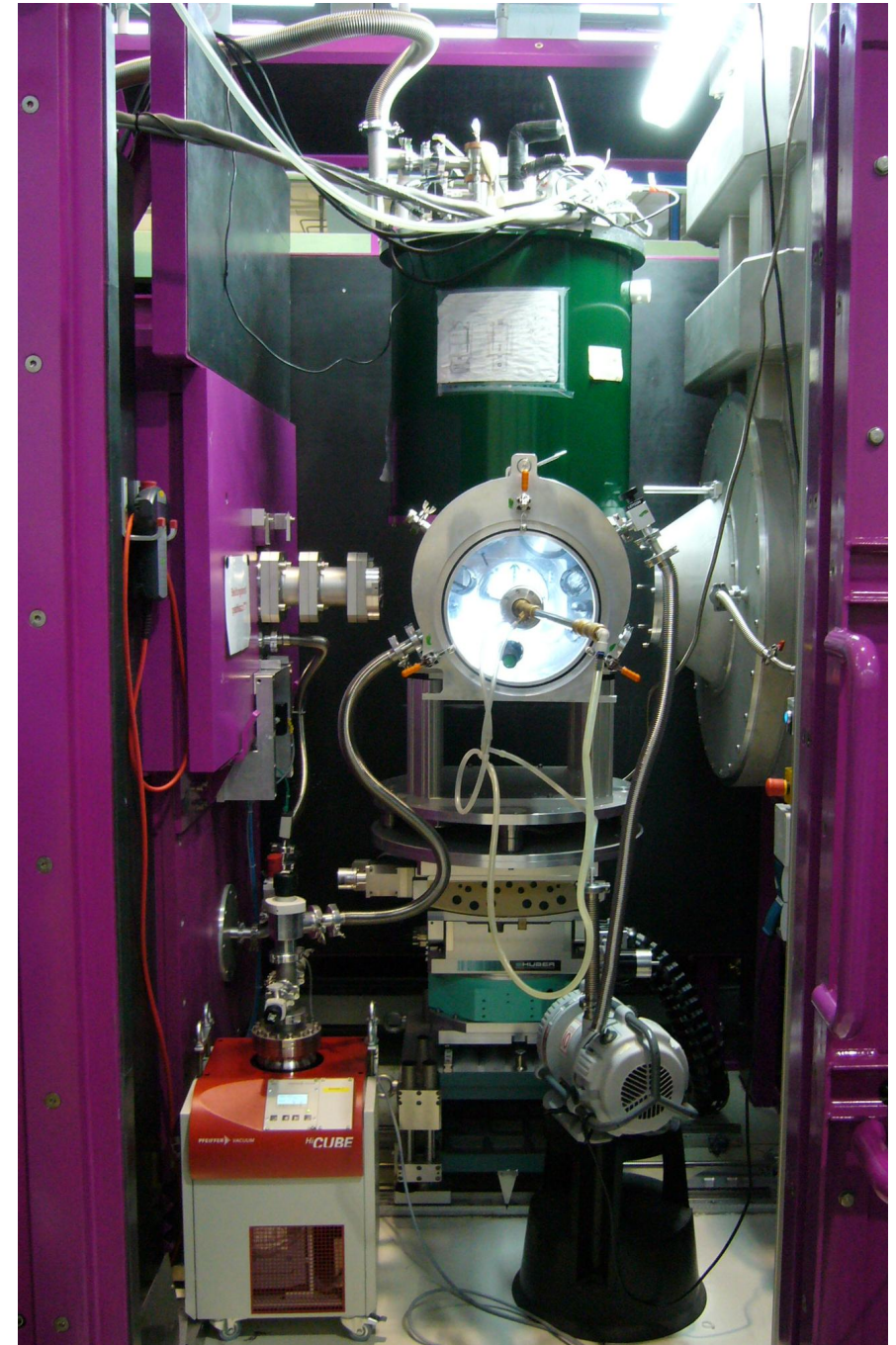
Parasitic background scattering has to be avoided (whole collimation system is evacuated, edge scattering, incoherent background)



Provide necessary sample environment (similar to any other neutron diffractometer or spectrometer)

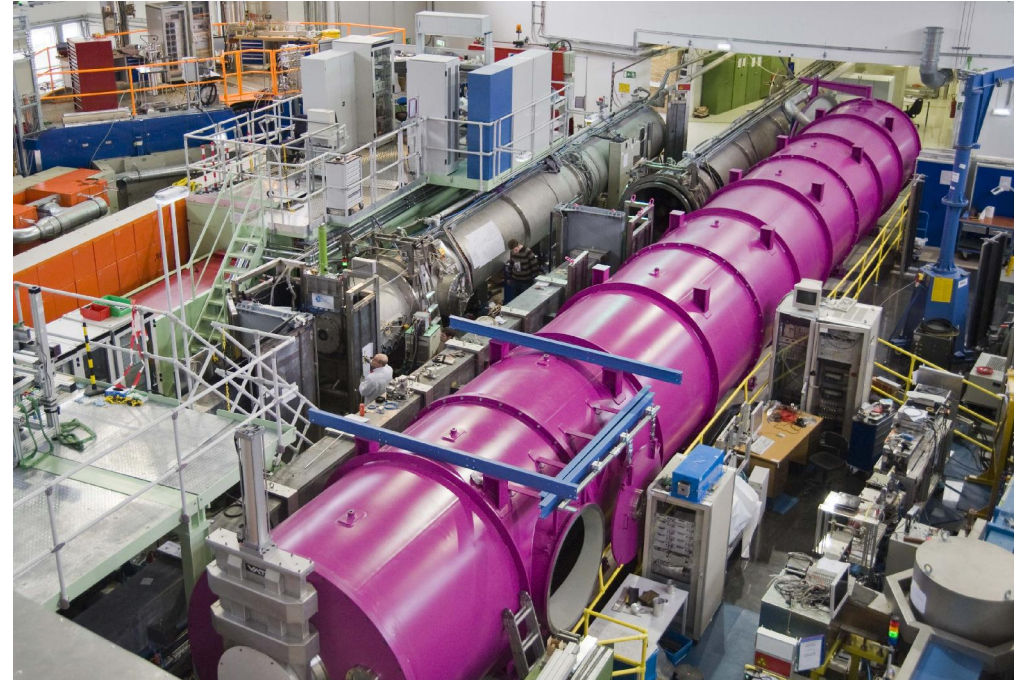
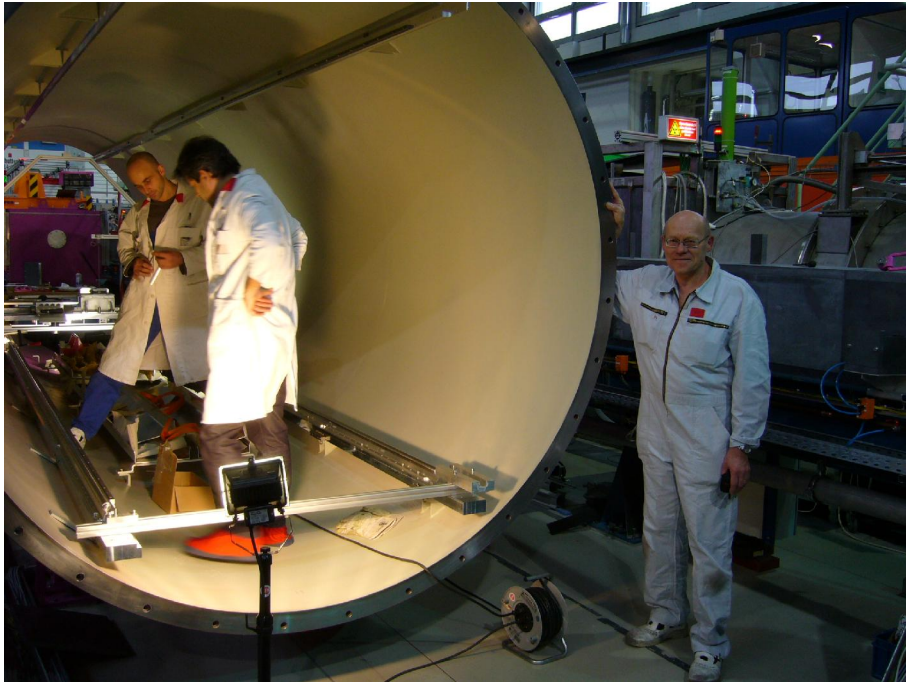
Parasitic background scattering has to be avoided (extremely critical!)

- ➡ Minimize neutrons travelling in air (few cm can be too much)
- ➡ Avoid Aluminum neutron windows (single crystalline sapphire is better)
- ➡ Get rid of scattering at edges (use conical slits)





Vacuum vessel for detector to provide lowest possible background



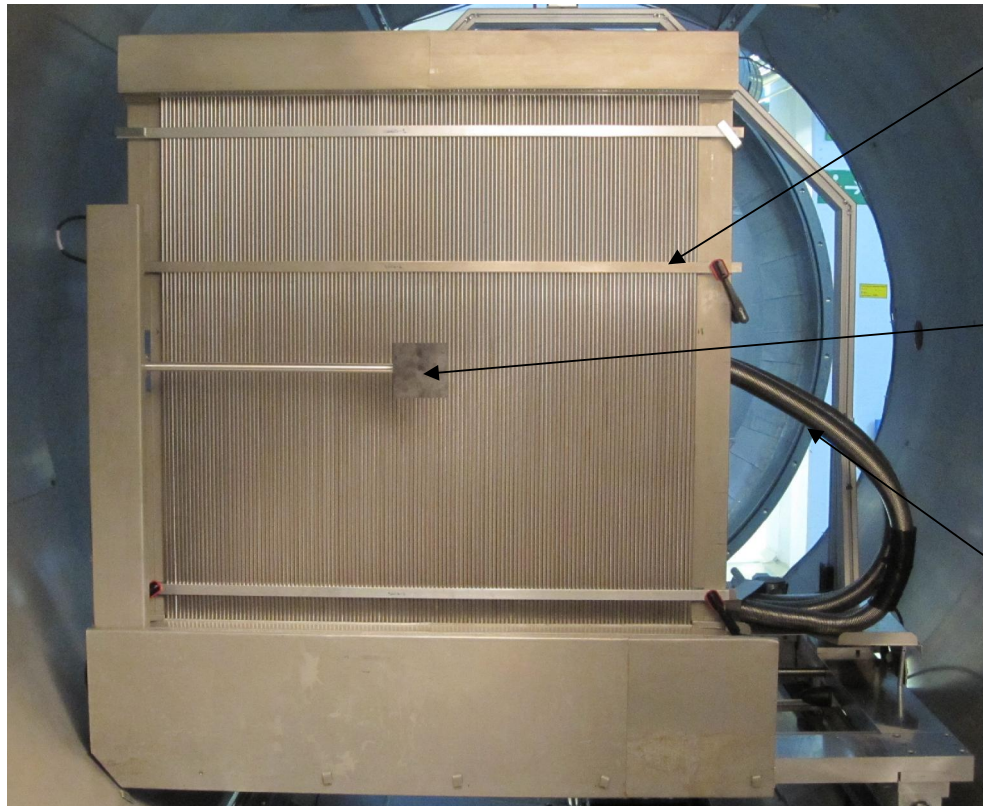
Sample detector length  
adjustable (select Q-range)

One (or several)  $\text{He}^3$  position  
sensitive detectors (typical  $1\text{m}^2$   
with 5mm resolution)

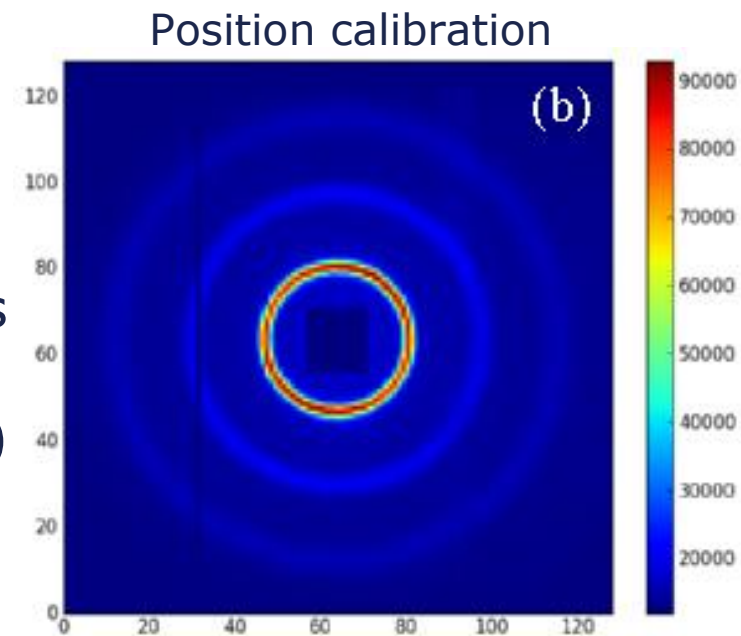
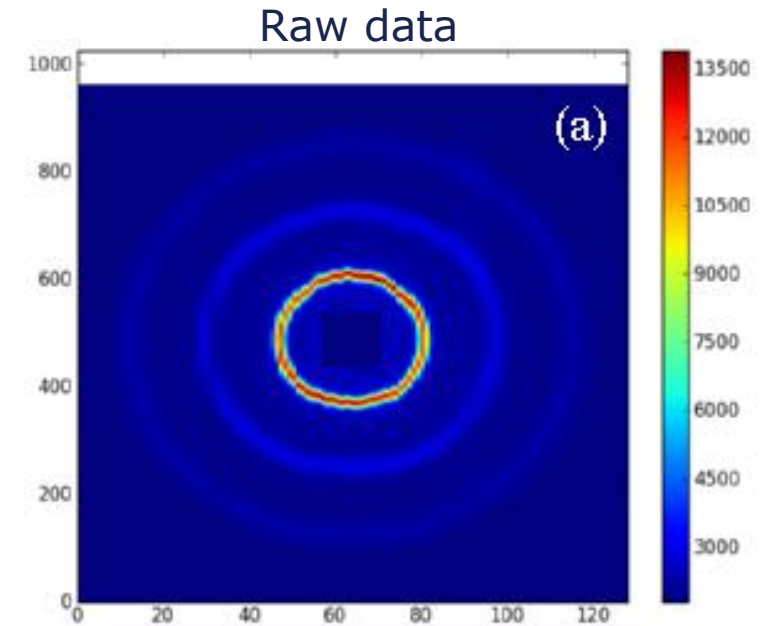
Typical length 10-40m

Interior completely covered with  
neutron absorbing Cadmium

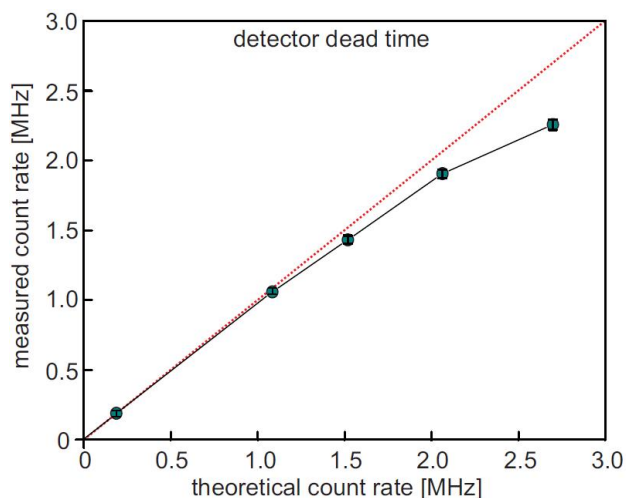




Array of 128 position sensitive  $^3\text{He}$  Reuter Stokes tube detectors  
8mm x 8mm position resolution (charge division)  
Detector distance from 1-20m, moves on rails  
Maximal count rate 4MHz

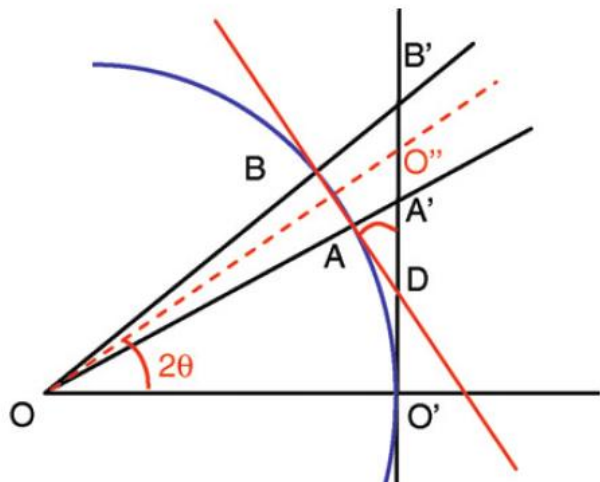


## Dead time corrections



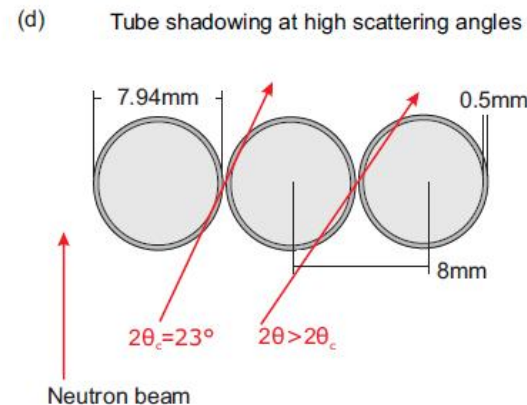
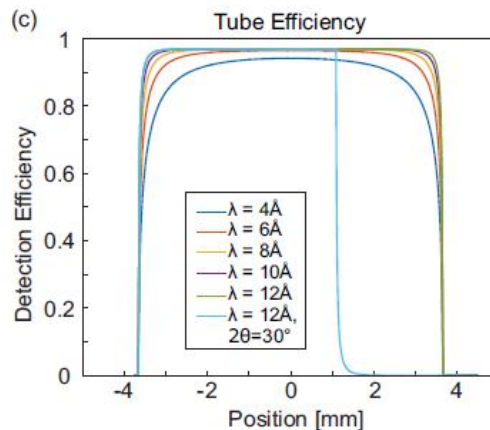
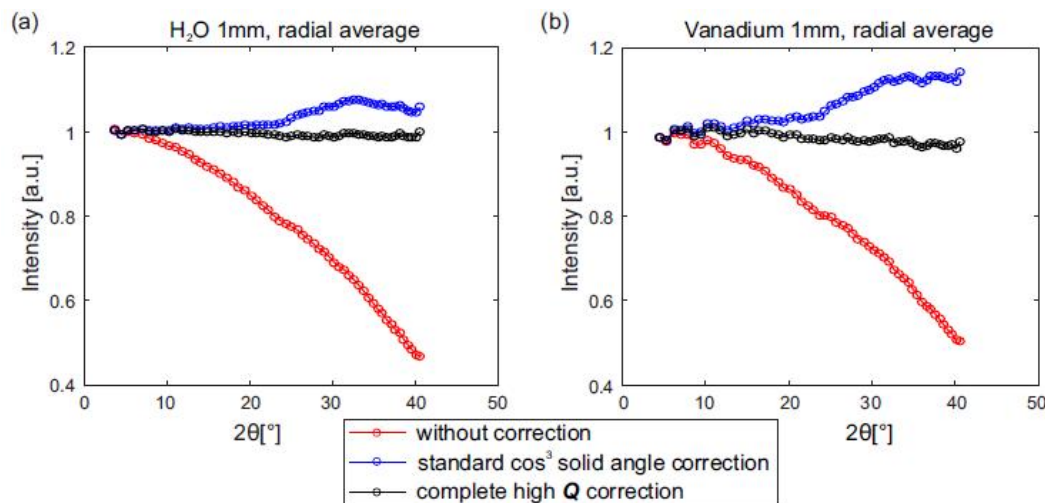
## Solid angle correction

$$\Delta\sigma(2\Theta) = \frac{p_x p_y \cos(2\Theta)^3}{D(2\Theta = 0)}$$

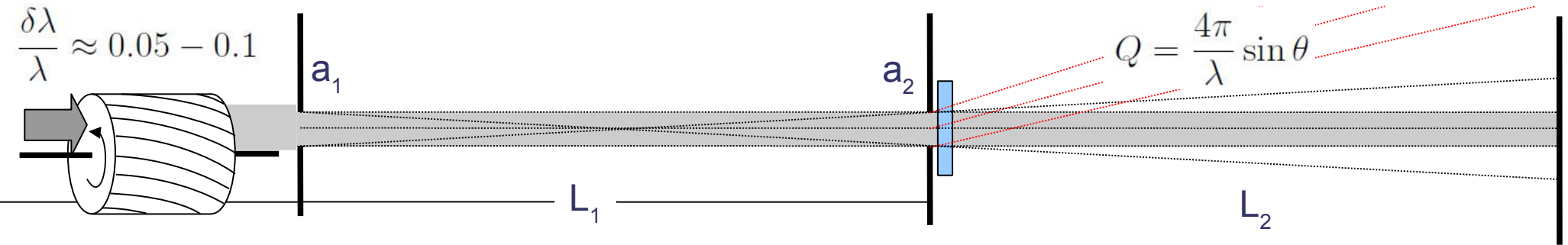


## Anisotropic solid angle correction for tube array detector

$$\Delta\sigma(2\Theta) = \frac{p_x p_y \cos(\Theta_x) \cos(2\Theta)^2}{D(2\Theta = 0)}$$



# SANS – Resolution & Intensity



Angular resolution  
Monochromaticity  
Detector resolution  
Gravity

Treat as Gaussian distributions:  $\left\langle \frac{\delta Q^2}{Q^2} \right\rangle = \left\langle \frac{\delta \lambda^2}{\lambda^2} \right\rangle + \left\langle \frac{\cos^2 \theta \delta \theta^2}{\sin^2 \theta} \right\rangle$

$$\left\langle \frac{\delta Q^2}{Q^2} \right\rangle = 0.0025 + \left\langle \frac{\delta \theta^2}{\theta^2} \right\rangle \Rightarrow \text{Angular resolution: } \delta \theta \approx \sqrt{\frac{5}{12} \frac{a}{L}}$$

What is the largest object SANS can detect  
(limit small Q)?

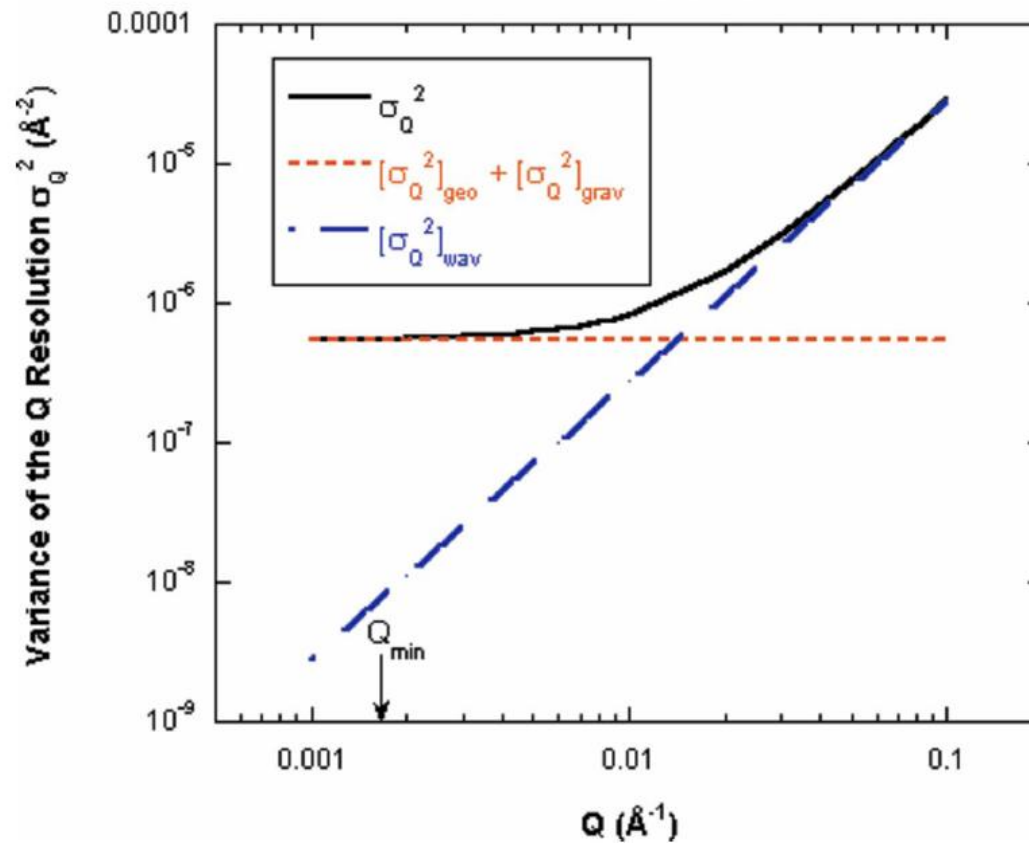
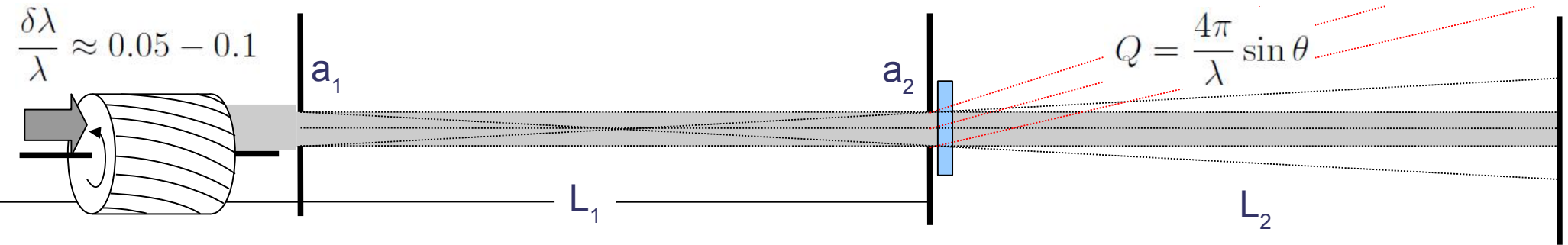
For large scattering angles (large Q)  
wavelength resolution dominates.

$$a_1 = a_2 = a \quad L_1 = L_2 = L$$

$$\Rightarrow \delta Q \approx \frac{\delta \theta}{\theta_{min}} Q_{min} \approx \delta \theta \frac{4\pi}{\lambda} \approx \frac{2\pi a}{\lambda L}$$

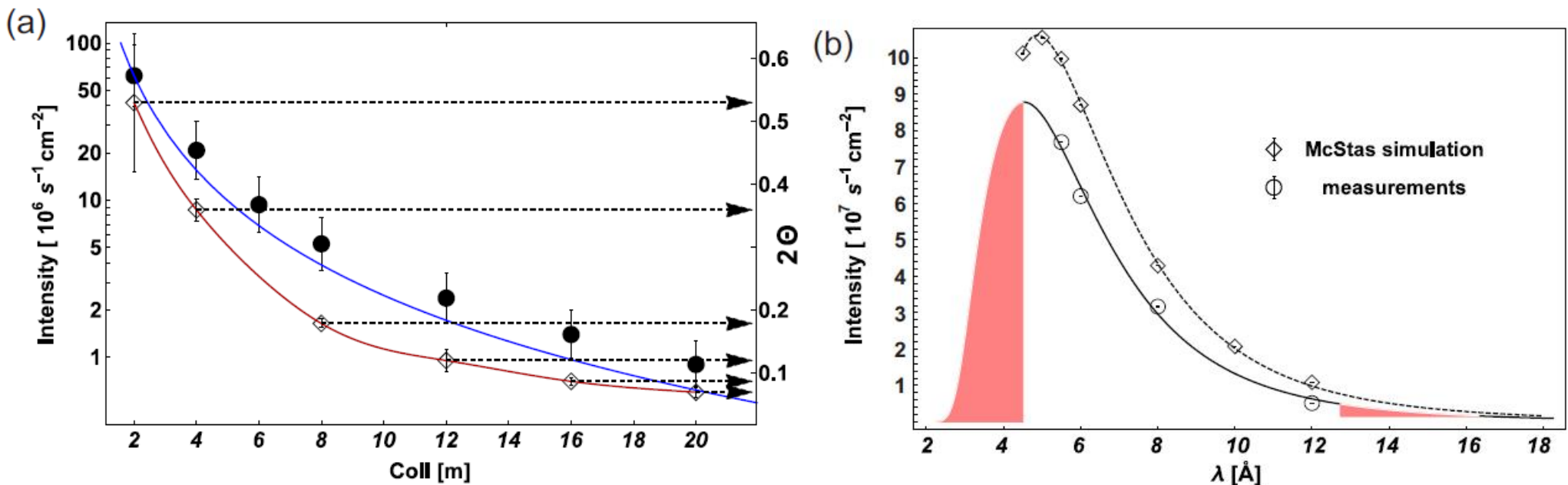
$$\frac{2\pi}{\delta Q} = \frac{\lambda L}{a}$$

On D11, ILL:  $L=40\text{m}$ ,  $\lambda=15\text{\AA}$   $\Rightarrow D \approx 5\mu\text{m}$





# SANS – Intensity & Resolution



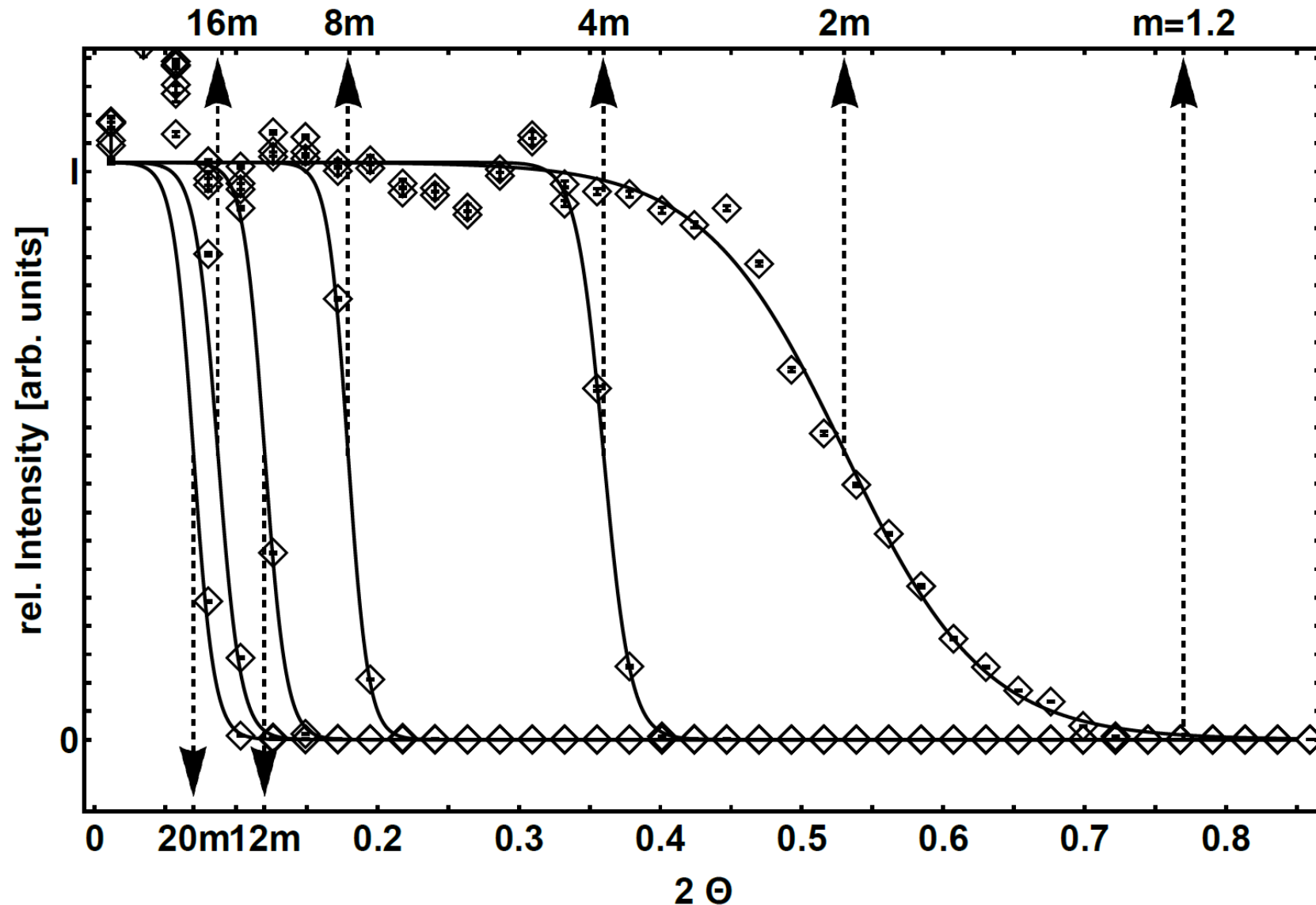
Intensity: Quadratic decrease  
with source to sample distance  
(collimation length)

Wavelength: Decrease of  
intensity with  $\lambda^{-4}$

$$\delta Q \approx \frac{\delta \theta}{\theta_{min}} Q_{min} \approx \delta \theta \frac{4\pi}{\lambda} \approx \frac{2\pi a}{\lambda L}$$



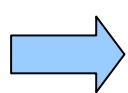
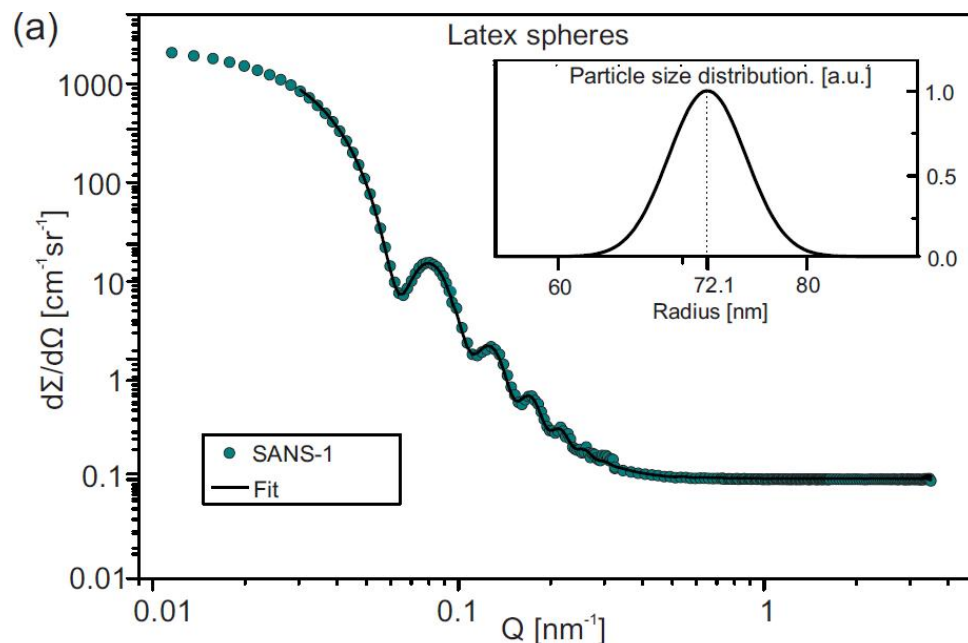
# SANS – Intensity & Resolution



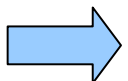
## Typical SANS dataset:

- Sample (at different L)
- Water (absolute scale)
- Empty sample holder/cuvette
- Background

$$\left(\frac{d\Sigma}{d\Omega}\right)_{\text{sample}} = \frac{1}{F_{\text{sc}}} \left(\frac{d\Sigma}{d\Omega}\right)_{\text{H}_2\text{O}}^{\text{real}} \frac{\left[\frac{I_{\text{sample}} - I_{\text{B4C}}}{Tr_{\text{sample}}} - \frac{I_{\text{sample-EC}} - I_{\text{B4C}}}{Tr_{\text{sample-EC}}}\right] \frac{1}{e_{\text{sample}}}}{\left[\frac{I_{\text{H}_2\text{O}} - I_{\text{B4C}}}{Tr_{\text{H}_2\text{O}}} - \frac{I_{\text{H}_2\text{O-EC}} - I_{\text{B4C}}}{Tr_{\text{H}_2\text{O-EC}}}\right] \frac{1}{e_{\text{H}_2\text{O}}}}$$



Fit model of the sample (conv. with resolution to the dataset)



SANS tells you:

- Shape of scattering object
- Size(distribution) of scattering objects
- Surface of scattering objects
- Scattering length density distribution
- Arrangement (Superstructure?)

# SANS – applications

Soft matter  
Hard matter  
Magnetism

## Polyisoprene – Polystyrene Diblock Copolymer Phase Diagram

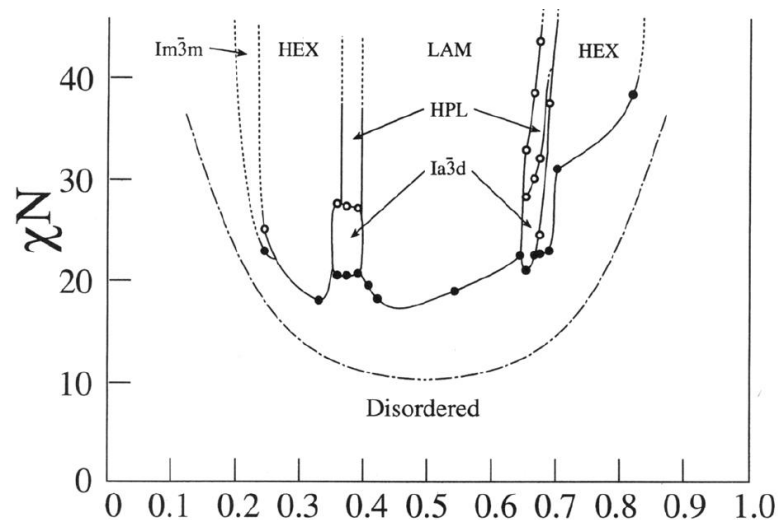
Two (or more) homopolymers units linked by covalent bonds

Microphase separation: Complex nanostructures phases

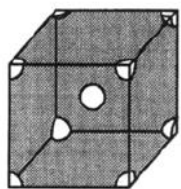
→ Flory-Huggins segment-segment interaction

→ Degree of polymerization

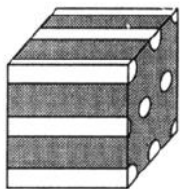
→ Volume fraction



$f_{PI}$



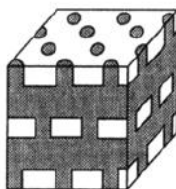
Spheres



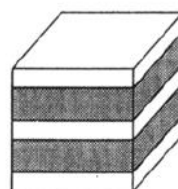
Cylinders



Bicontinuous

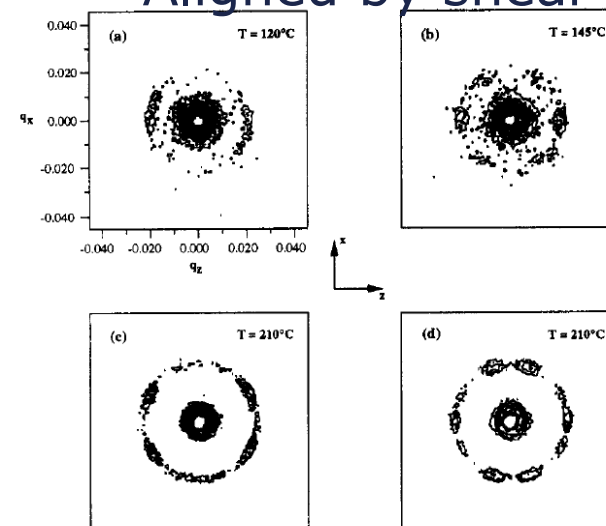


Perforated  
Layers



Lamellae

SANS pattern:  
Aligned by shear



**Figure 10.** Contour plots of SANS patterns for the sample IS-68 in three different ordered states (A, D, B). A nearly featureless pattern (a) was observed in state A (120 °C), which is attributed to a parallel orientation of lamellae with respect to the shear plane. After heating the sample to 145 °C (state D) and application of dynamic shearing ( $\dot{\gamma} = 0.1 \text{ s}^{-1}$  with  $|\gamma| = 300\%$ ), a weak hexagonal scattering pattern was observed (b). This is consistent with a hexagonal in-plane arrangement of perforations in the minority (PS) layers. Further heating the sample to 210 °C, without shear (state B), produced a predominantly four-peak pattern (c), which transformed into result (d) when a shear rate of  $2.2 \text{ s}^{-1}$  was applied. The azimuthal relationship between the combined 10 reflections in (c) and (d) is consistent with the SANS data reported for sample IS-39 (ref 22) and the  $Ia\bar{3}d$  space group symmetry.

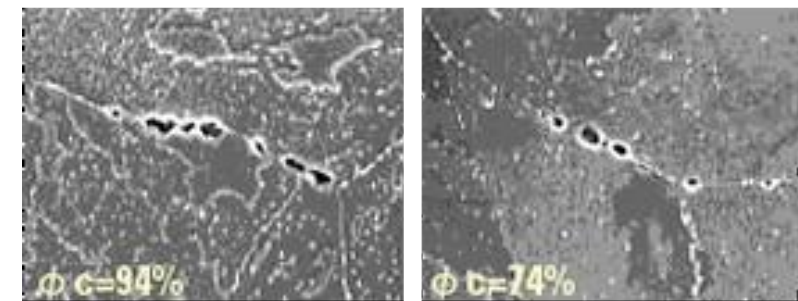
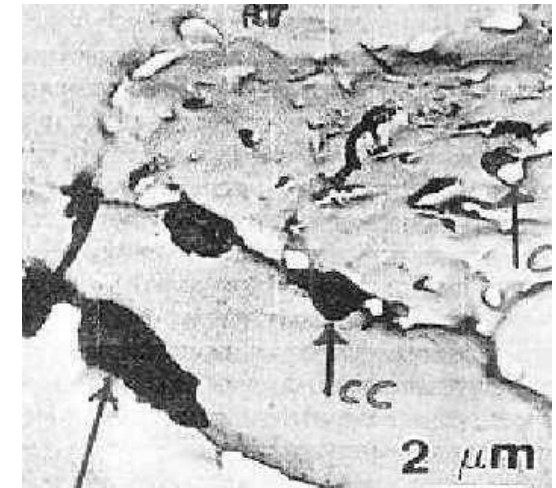
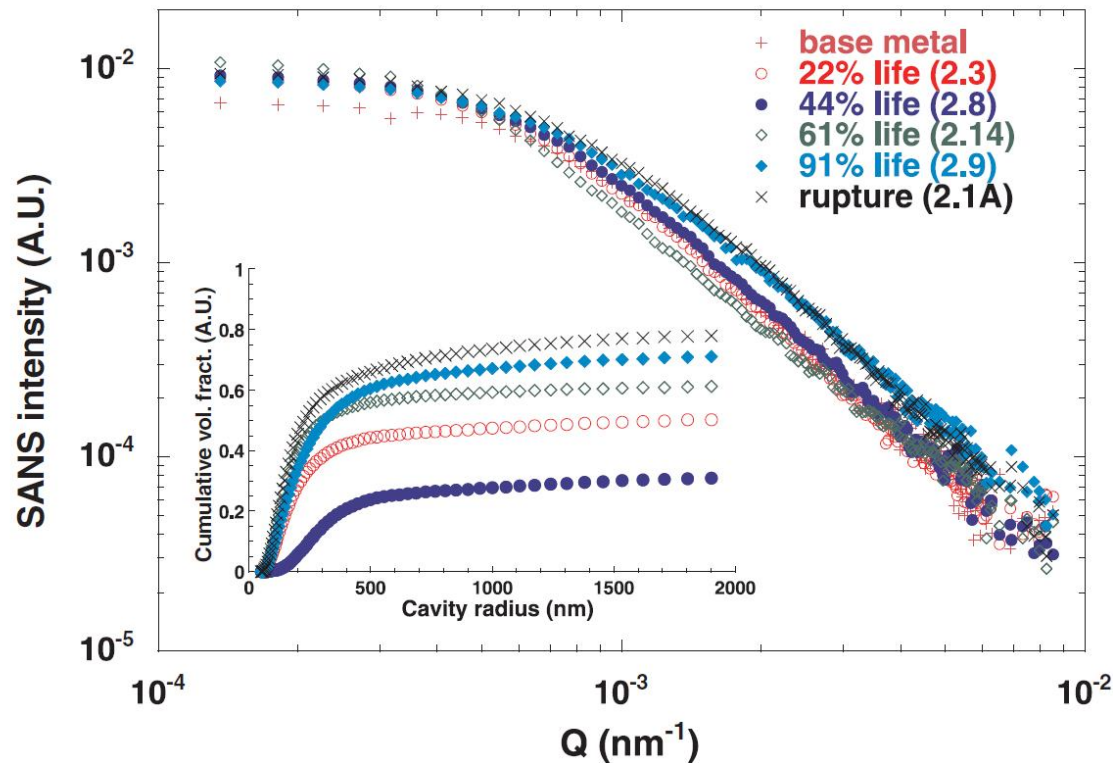
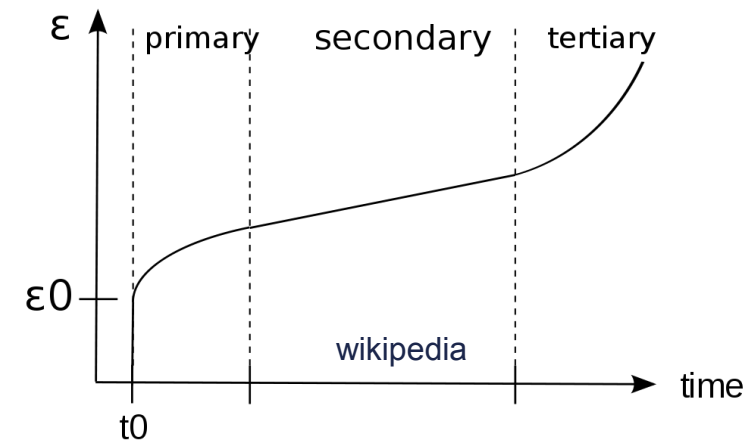
Macromolecules, Vol. 28, No. 26,  
(1995)

Physics Today, p. 32, Feb. (1999)

## Creep cavitation damage in steel at high T

Volume fraction and size distribution of cavities can be measured with SANS and USANS.

Grain boundary cavitation dominant failure mode.

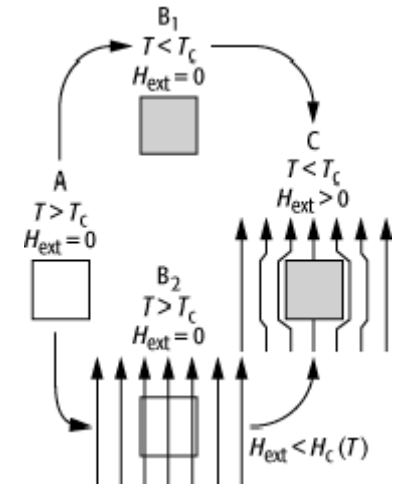
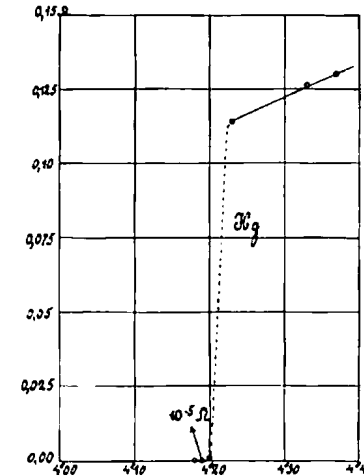


**Fig. 3.** SANS intensity for the investigated specimens, measured in the double-crystal experiment at HMI-BENSC (Inset: Cumulative cavity volume fraction as a function of cavity radius)



## Superconductivity (H. Kamerlingh Onnes 1911)

- ➡ Total loss of resistivity
- ➡ Expulsion of magnetic fields (Meissner Ochsensfeld Effect)



Two length scales:

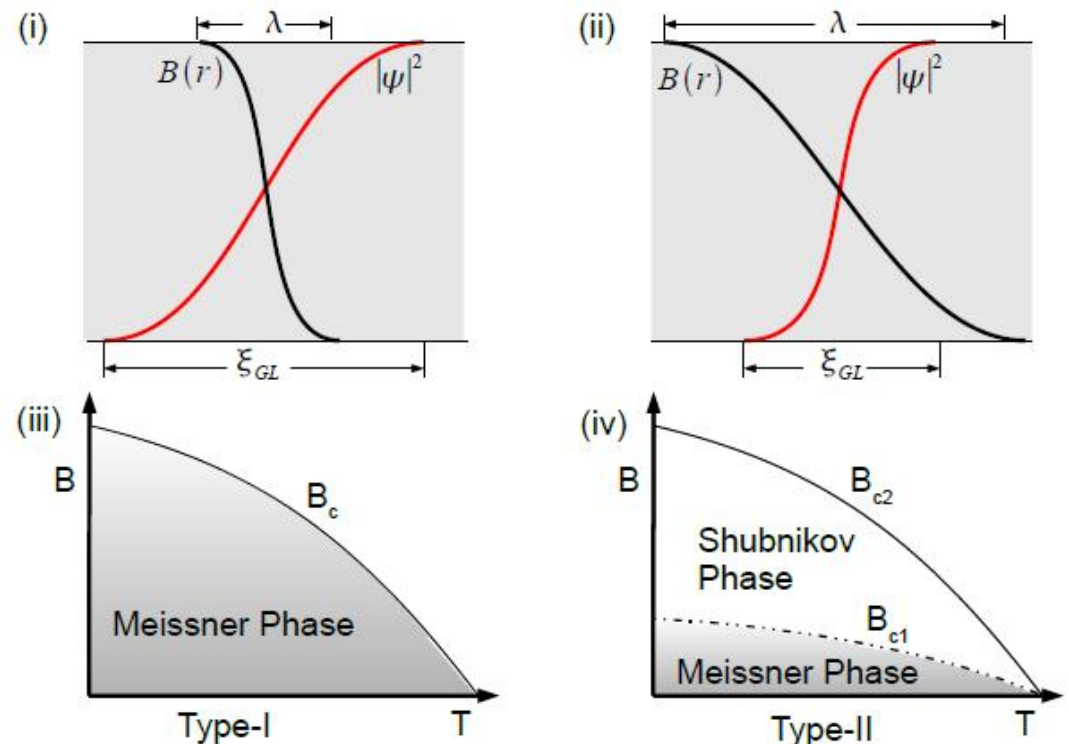
Penetration depth

$$\lambda^2(T) = \frac{m^* c^2}{4\pi n_s e^{*2}}$$

Coherence length

$$\xi^2(T) = \frac{\hbar^2}{2m^* |\alpha(T)|} \propto \frac{1}{1 - \frac{T}{T_c}}$$

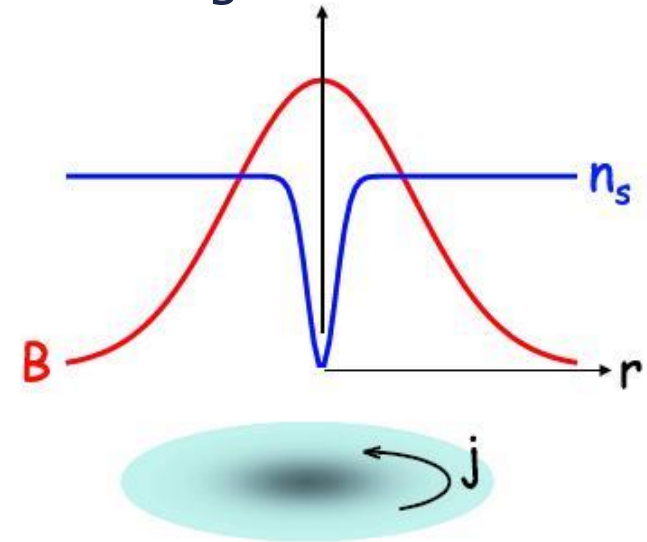
- ➡ Interface energy





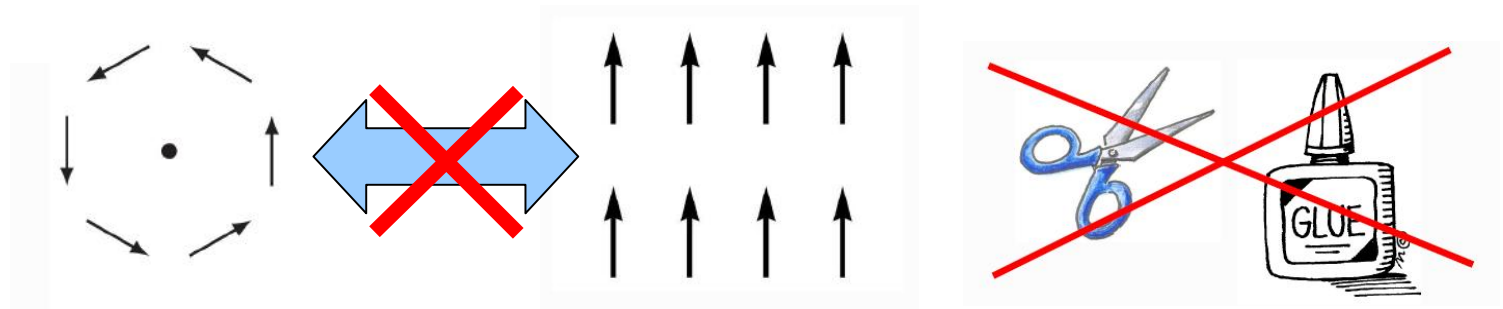
Stabilized by negative energy of super-/normal conducting interface.

➔ Quantization of magnetic flux  $\phi_0 = \frac{h}{2e}$



Now consider the topology!

Superconducting vortex: Topological defect of the superconducting OP.  
No continuous transformation from no vortex to a vortex state.

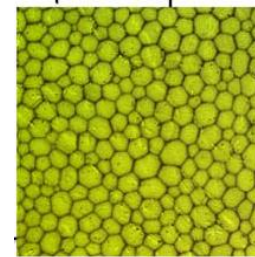
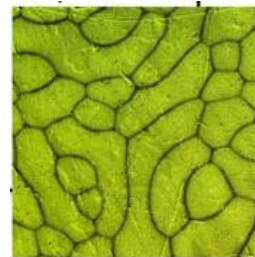


Protected by topology: Particle-like properties

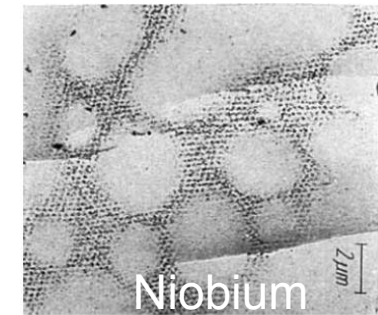
Condensed matter  $\longleftrightarrow$  superconducting vortex matter

Properties of superconducting VM reflect underlying physics  
Model system for general questions

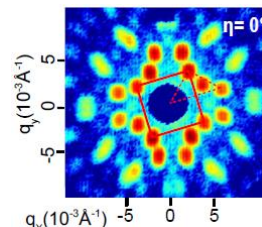
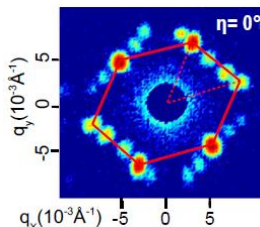
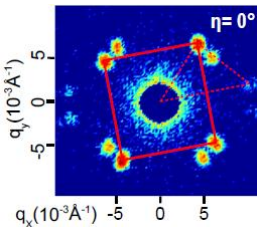
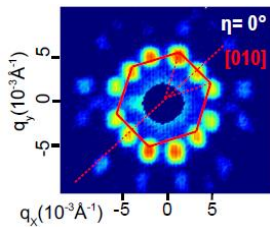
Domain structure



Lead

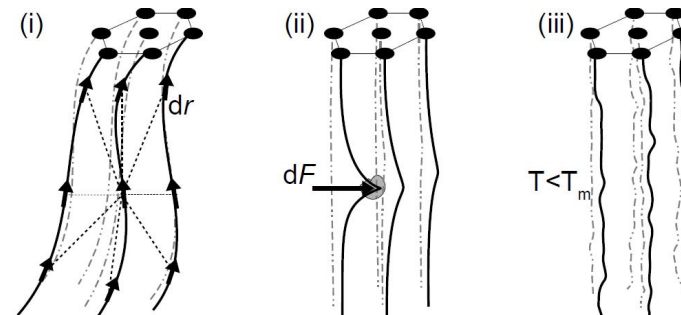


Niobium



Symmetry and structure

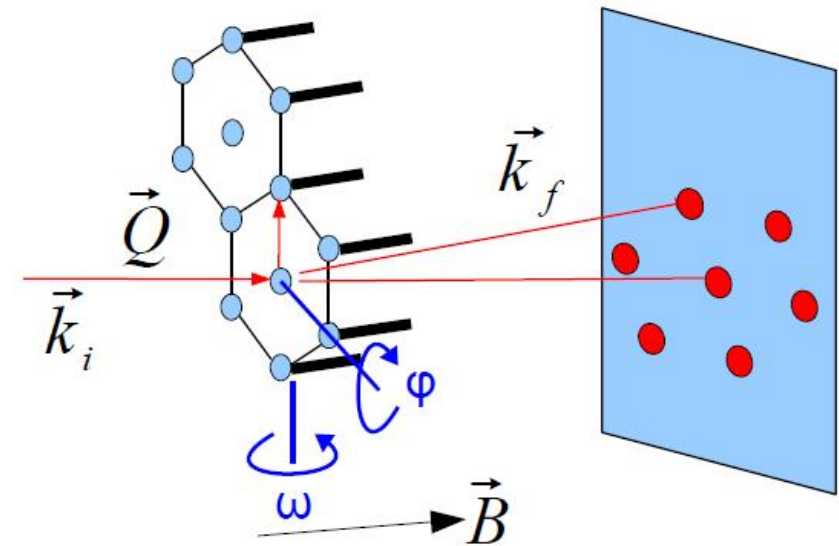
Elasticity & melting



## Vortex lattice 2D magnetic Bravais lattice

➔ One flux quantum per unit cell

$$\phi_0 = \frac{h}{2e} \quad |\vec{a}_i| = \left( \frac{2\phi_0}{\sqrt{3}B} \right)^{\frac{1}{2}} \quad |\vec{a}_i| = \frac{2\pi}{|\vec{Q}|}$$



➔ Typical values:

$$B = 1500 \text{ G}$$

$$A_0 = 1260 \text{ \AA}$$

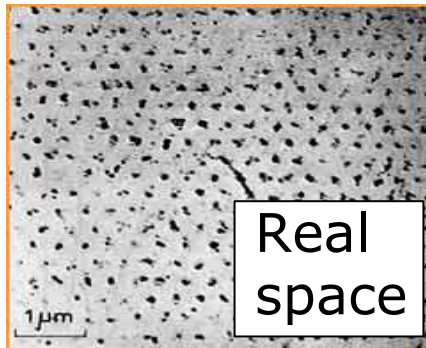
$$|\vec{Q}| = 0.0057 \text{ \AA}^{-1}$$

Intensity Bragg peak

$$R = \frac{2\pi \gamma^2 \lambda_n^2 t}{16\phi_0^2 Q} |h(Q)|^2$$

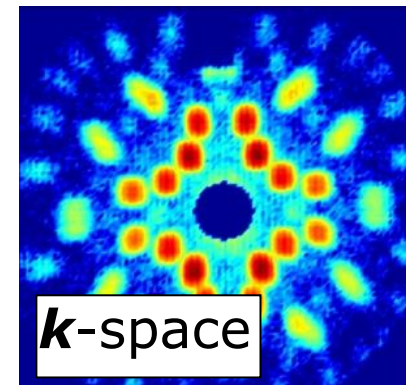
Form factor

$$h(Q) = \frac{\phi_0}{(2\pi\lambda)^2} e^{\frac{-\pi B}{B_{c2}}}$$



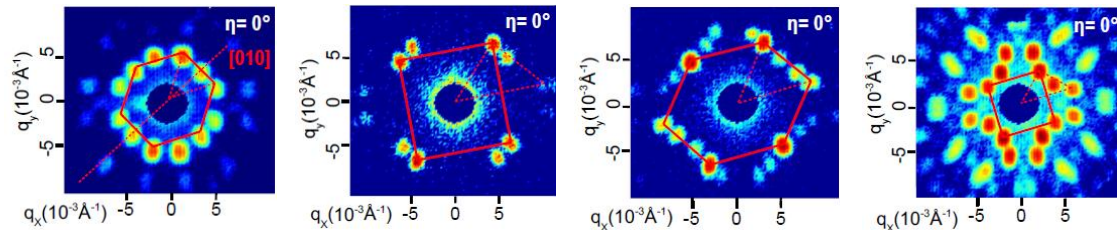
- Rocking gives all bragg peaks

90° rot. around the n-beam



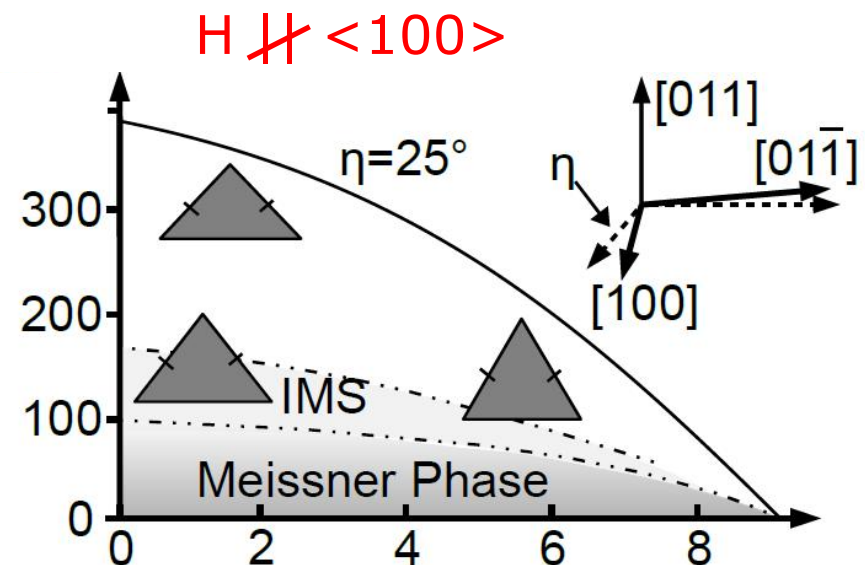
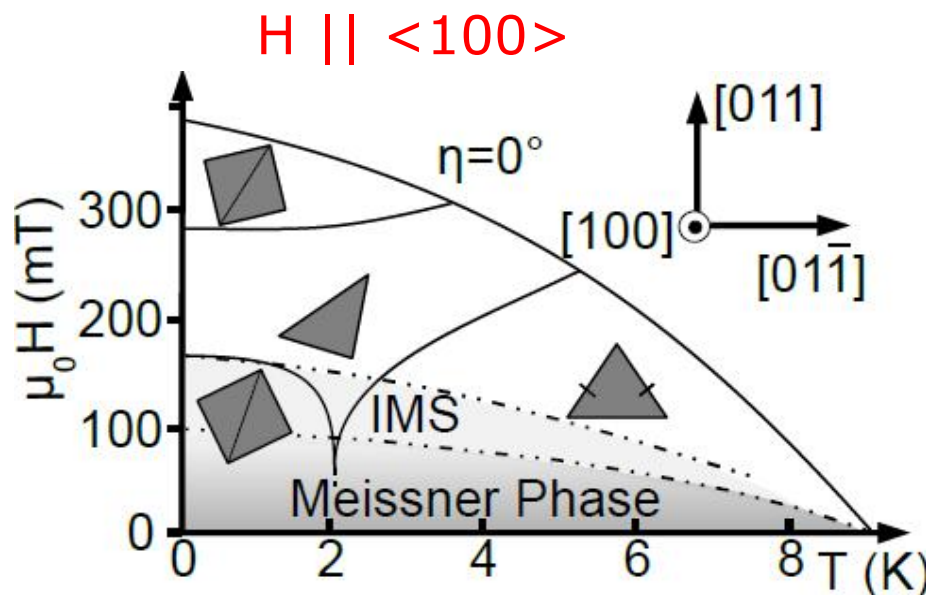
## Symmetry and structure

Nature of OP, symmetry of Fermi surface. Intricate to separate different influences.



S. Mühlbauer et al. *Phys. Rev. Lett.* 102, 136408 (2009)

Six-fold symmetry of vortex lattice  $\longleftrightarrow$  Frustration  $\longleftrightarrow$  Four fold crystal symmetry  $\rightarrow$  Four VL phases, break crystal symmetry!

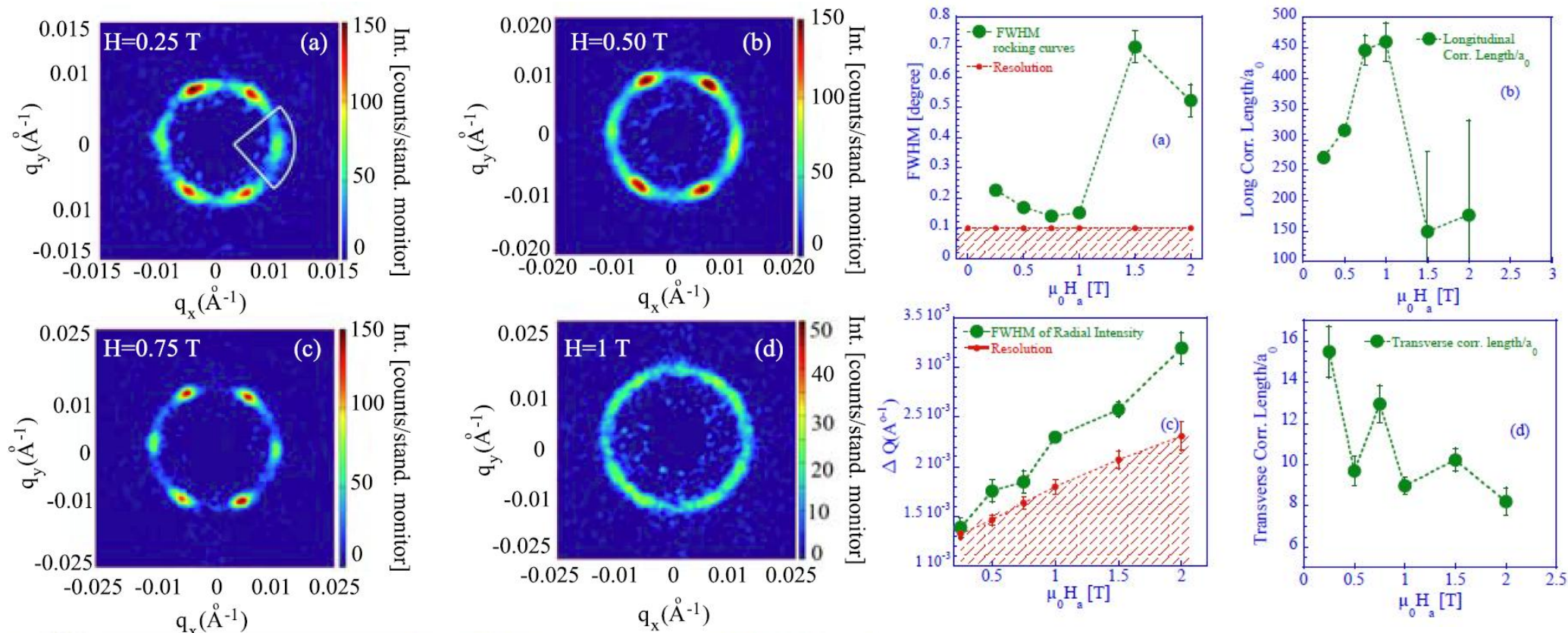




## Structure & form factor, correlation lengths

### Optimally doped $\text{Ba}_{(1-x)}\text{K}_x\text{Fe}_2\text{As}_2$

#### Longitudinal and transverse correlation lengths of the vortex lattice



S. Demirdis et al. submitted to *PRB* (2015)

# SANS – extensions to lower Q

Focusing  
VSANS

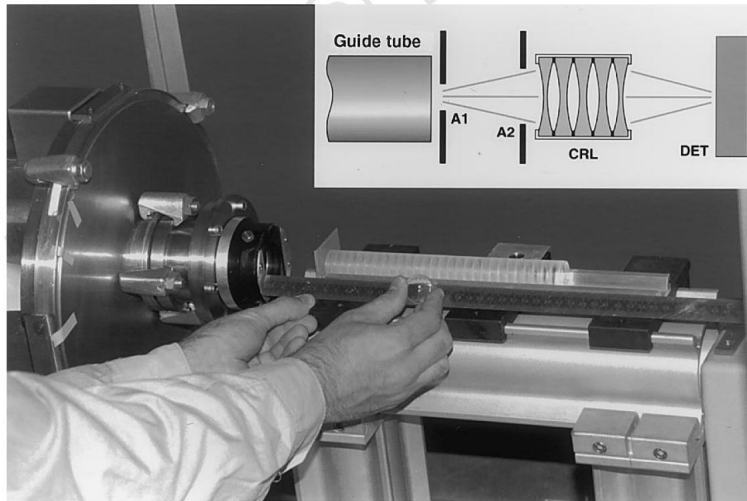
USANS (Bonse Hart Camera)  
MSANS



Refractive index  $< 1$   $n = 1 - \rho b_c \frac{\lambda^2}{2\pi}$

Focal length (single  $\text{MgF}_2$  lens)  $\approx 200\text{m}$

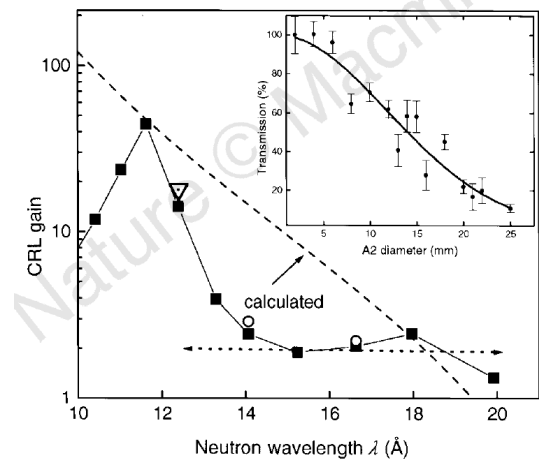
→ Stack of lenses is used



Nature **391**, p 563 (1998)

Ideal lens:

- No absorption
- No incoherent scattering
- Large refractive index

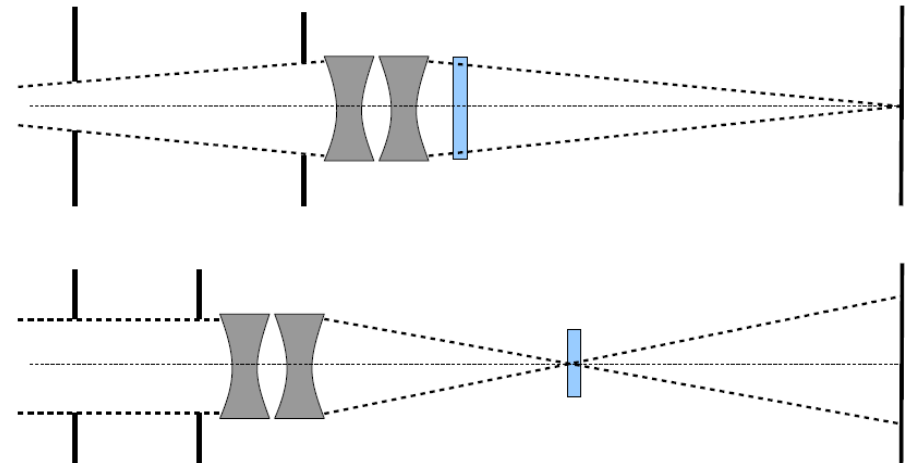


Boosting the resolution: Focus the neutron beam on the detector.

→ Large sample needed.

Boosting the intensity: Focus the neutron beam on the sample.

→ Sacrifice Q-resolution.



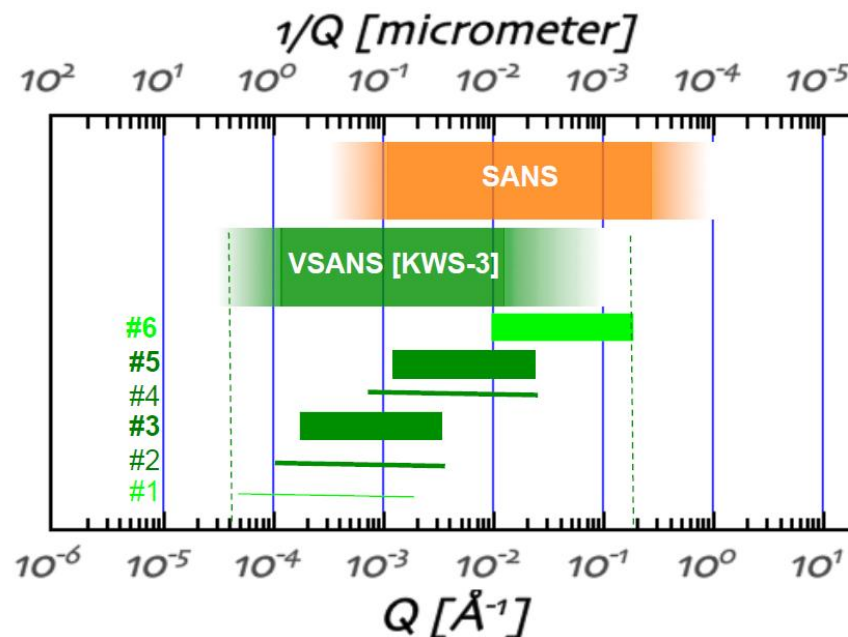
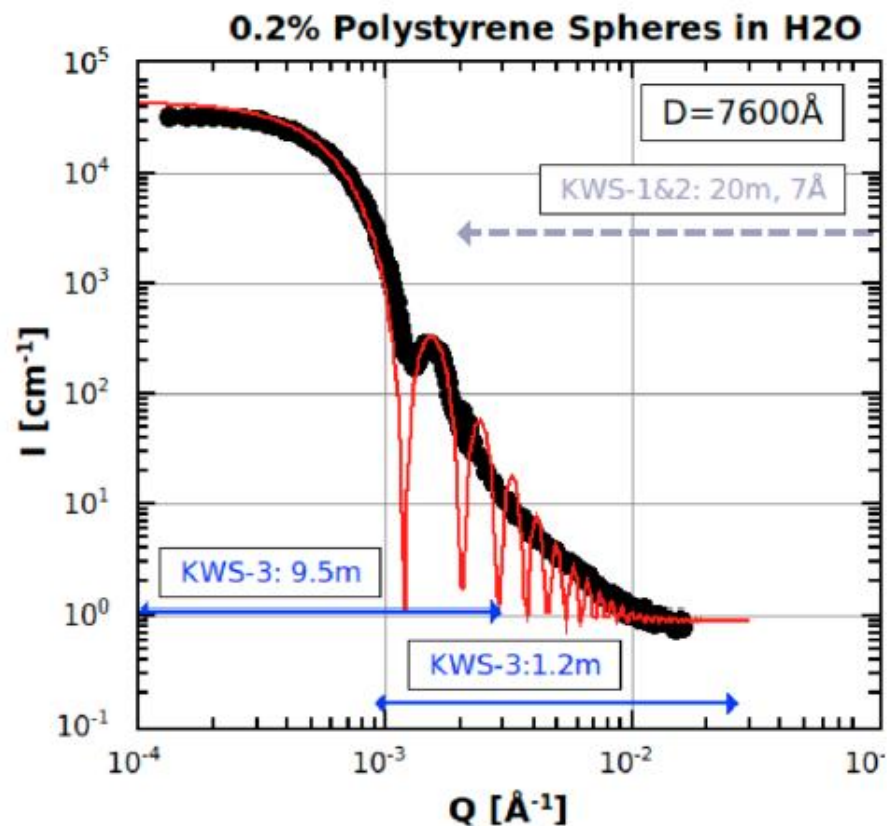
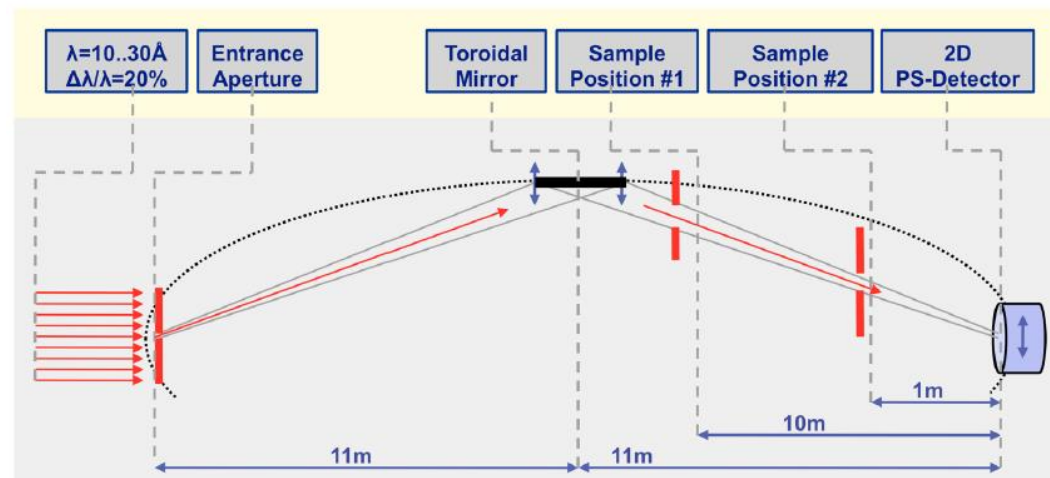
## VSANS

(Toroidal elliptical mirror)

$$10^{-4} \text{ \AA}^{-1} < Q < 3 \cdot 10^{-3} \text{ \AA}^{-1}$$

KWS 3 (FRM II)

Instrument layout

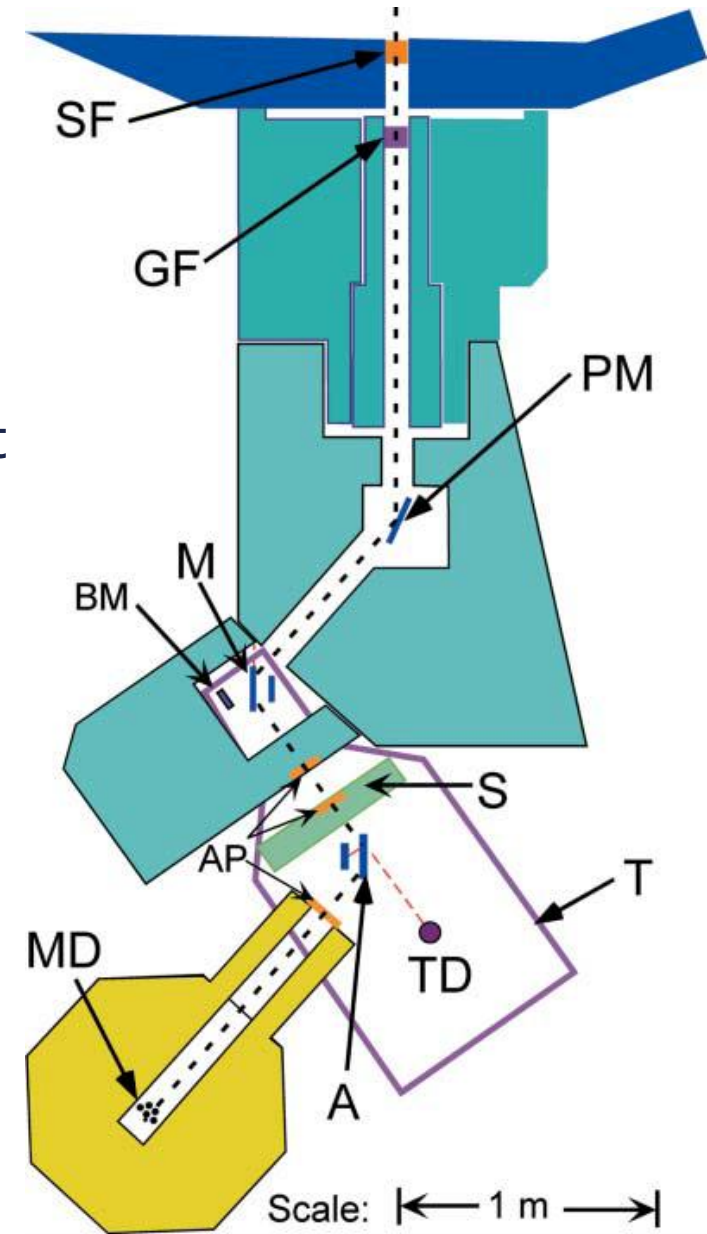
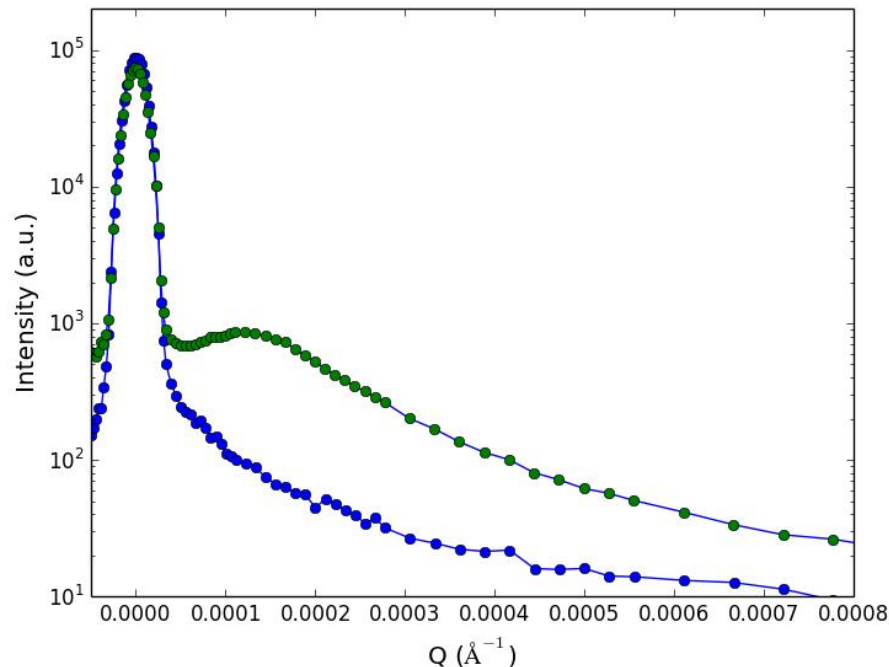


## USANS (Bonse-Hart Camera)

$$3 \cdot 10^{-5} \text{ \AA}^{-1} < Q < 5 \cdot 10^{-3} \text{ \AA}^{-1}$$

S18 (ILL), BT5 (NIST), Kookaburra (ANSTO), V12 (HMI)

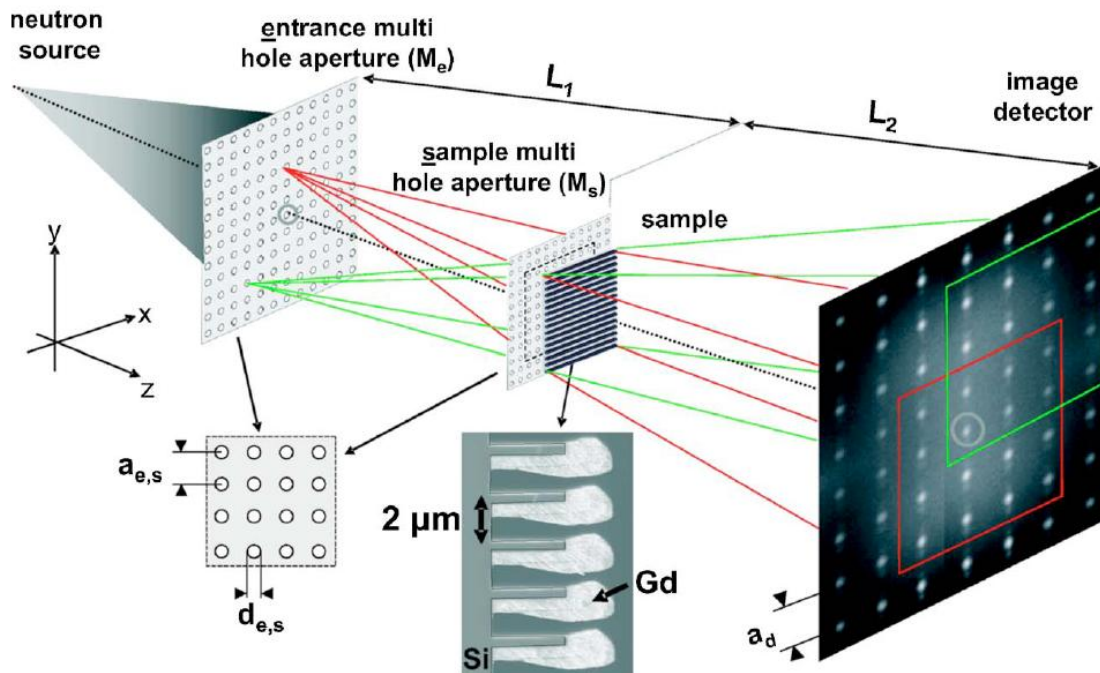
- ➡ Use multiple reflection at channel cut perfect single crystal monochromators
- ➡ Slit smeared data



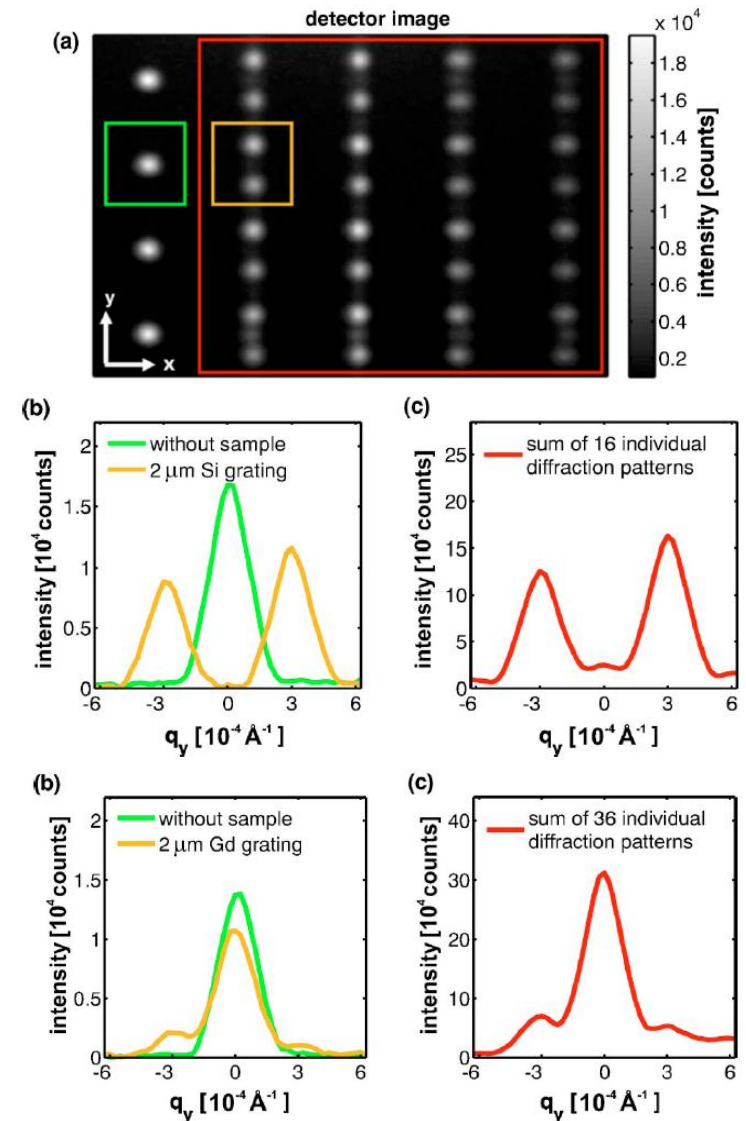
Multi-pinhole masks.

Theorem of intersecting lines.

Coherent summation of different scattering patterns at the detector



$$\frac{a_d}{a_s} = \frac{L_1 + L_2}{L_1} \quad \frac{a_d}{a_e} = \frac{L_2}{L_1}$$



Resolution of  $3 \cdot 10^{-4} \text{Å}^{-1}$  with 2.6m collimation  
Useful for particular samples and Q-range  
Problem: Edge scattering and background!