

Notes for exercise sheet 8

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Exercise 1

As given in a previous exercise, the bcc lattice has 8 nearest neighbours at the positions:

$$\vec{n}_1 = \begin{pmatrix} 1/2 \\ 1/2 \\ 1/2 \end{pmatrix}, \vec{n}_2 = \begin{pmatrix} -1/2 \\ 1/2 \\ 1/2 \end{pmatrix}, \vec{n}_3 = \begin{pmatrix} 1/2 \\ -1/2 \\ 1/2 \end{pmatrix}, \vec{n}_4 = \begin{pmatrix} 1/2 \\ 1/2 \\ -1/2 \end{pmatrix} \quad (1)$$

$$\vec{n}_5 = \begin{pmatrix} -1/2 \\ -1/2 \\ 1/2 \end{pmatrix}, \vec{n}_6 = \begin{pmatrix} -1/2 \\ 1/2 \\ -1/2 \end{pmatrix}, \vec{n}_7 = \begin{pmatrix} 1/2 \\ -1/2 \\ -1/2 \end{pmatrix}, \vec{n}_8 = \begin{pmatrix} -1/2 \\ -1/2 \\ -1/2 \end{pmatrix}. \quad (2)$$

The antiferromagnetic dispersion relation is given by:

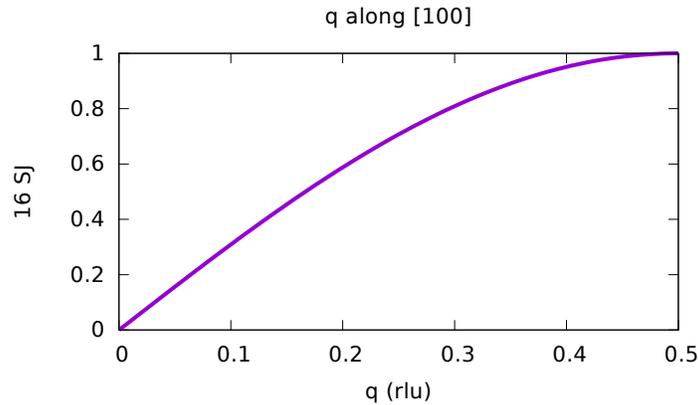
$$E = 2S \cdot \left[\left(\sum_{j \in \{1, \dots, 8\}} J_j \right)^2 - \left(\sum_{j \in \{1, \dots, 8\}} J_j \exp(-i2\pi\vec{q} \cdot \vec{n}_j) \right)^2 \right]^{1/2} \quad (3)$$

$$E = 2S \cdot \left[64J^2 - \left(\sum_{j \in \{1, \dots, 8\}} J_j \exp(-i2\pi\vec{q} \cdot \vec{n}_j) \right)^2 \right]^{1/2} \quad (4)$$

- Along an [100] direction:

$$E_{[100]} = 2S \cdot \left[64J^2 - (4J \exp(-i\pi q_x) + 4J \exp(i\pi q_x))^2 \right]^{1/2} \quad (5)$$

$$E_{[100]} = 2S \cdot [64J^2 - 64J^2 \cos^2(\pi q_x)]^{1/2} = 16SJ \cdot [1 - \cos^2(\pi q_x)]^{1/2} \quad (6)$$



- Along an [110] direction:

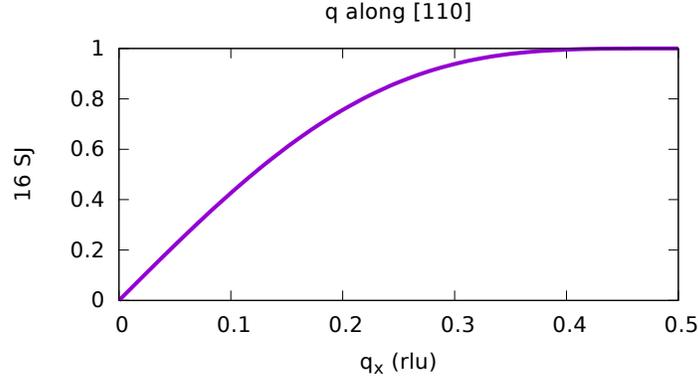
$$E_{[110]} = 2S \cdot \left[64J^2 - \left(2J \exp(i\pi q_x + i\pi q_y) + 2J \exp(-i\pi q_x - i\pi q_y) + \overbrace{2J \exp(-i\pi q_x + i\pi q_y)}^{=2J} + \overbrace{2J \exp(i\pi q_x - i\pi q_y)}^{=2J} \right) \right]^{1/2} \quad (7)$$

Using $q_x = q_y$:

$$E_{[110]} = 2S \cdot \left[64J^2 - (4J + 2J \exp(i2\pi q_x) + 2J \exp(-i2\pi q_x))^2 \right]^{1/2} \quad (8)$$

$$E_{[110]} = 2S \cdot \left[64J^2 - (4J(1 + \cos(2\pi q_x)))^2 \right]^{1/2} \quad (9)$$

$$E_{[110]} = 8SJ \cdot \left[4 - (1 + \cos(2\pi q_x))^2 \right]^{1/2} \quad (10)$$



Exercise 2

$$\mathcal{H} = -J\vec{S}_1 \cdot \vec{S}_2 = -J(S_{1,x}S_{2,x} + S_{1,y}S_{2,y} + S_{1,z}S_{2,z}) \quad (11)$$

Using the ladder operators $S^\pm = S_x \pm iS_y$:

$$\mathcal{H} = -\frac{1}{2}J(S_1^+S_2^- + S_1^-S_2^+) - JS_{1,z}S_{2,z} \quad (12)$$

Applying the Hamilton to the four possible 2-spin states:

$$\mathcal{H}|\uparrow\uparrow\rangle = -\frac{1}{2}J(0+0) - J\frac{1}{2}\frac{1}{2}|\uparrow\uparrow\rangle = -\frac{1}{4}|\uparrow\uparrow\rangle \quad (13)$$

$$\mathcal{H}|\downarrow\downarrow\rangle = -\frac{1}{2}J(0+0) - J\frac{1}{2}\frac{1}{2}|\downarrow\downarrow\rangle = -\frac{1}{4}|\downarrow\downarrow\rangle \quad (14)$$

$$\mathcal{H}|\uparrow\downarrow\rangle = -\frac{1}{2}J(|\uparrow\downarrow\rangle + 0) + J\frac{1}{2}\frac{1}{2}|\uparrow\downarrow\rangle = -\frac{1}{2}J|\uparrow\downarrow\rangle + \frac{1}{4}J|\uparrow\downarrow\rangle \quad (15)$$

$$\mathcal{H}|\downarrow\uparrow\rangle = -\frac{1}{2}J(0 + |\downarrow\uparrow\rangle) + J\frac{1}{2}\frac{1}{2}|\downarrow\uparrow\rangle = -\frac{1}{2}J|\downarrow\uparrow\rangle + \frac{1}{4}J|\downarrow\uparrow\rangle \quad (16)$$

The mixed eigenstates are found to be:

- triplet state:

$$|t\rangle = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle) \quad (17)$$

$$\mathcal{H}|t\rangle = \frac{1}{\sqrt{2}} \left(-\frac{J}{4} |\uparrow\downarrow\rangle - \frac{J}{4} |\downarrow\uparrow\rangle \right) \quad \text{eigenvalue : } -\frac{J}{4} \quad (18)$$

- singlet state:

$$|s\rangle = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \quad (19)$$

$$\mathcal{H}|s\rangle = \frac{1}{\sqrt{2}} \left(\frac{3}{4}J |\uparrow\downarrow\rangle - \frac{3}{4}J |\downarrow\uparrow\rangle \right) \quad \text{eigenvalue : } \frac{3}{4}J \quad (20)$$